Transverse Centroid and Envelope Descriptions of Beam Evolution*

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Transverse Centroid and Envelope Model: Outline

Comments:

- Some of this material related to J.J. Barnard lectures:
  - Transport limit discussions (Introduction)
  - Transverse envelope modes (Continuous Focusing Envelope Modes and Halo)
  - Longitudinal envelope evolution (Longitudinal Beam Physics III)
  - 3D Envelope Modes in a Bunched Beam (Cont. Focusing Envelope Modes and Halo)
- Specific transverse topics will be covered in more detail here for *s*-varying focusing
- Extensive Review paper covers envelope mode topics presented in more detail:
  Lund and Bukh, "Stability properties of the transverse envelope equations describing intense ion beam transport," PRSTAB **7** 024801 (2004)

Transverse Centroid and Envelope Model: Detailed Outline

Detailed Outline - 2

Detailed Outline - 3

10) Centroid and Envelope Descriptions via 1st Order Coupled Moment Equations
- Formulation
- Example Illustration -- Familiar KV Envelope Model

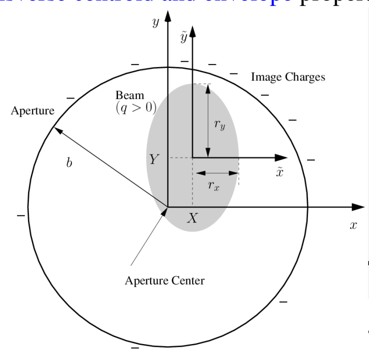
Contact Information

References

Acknowledgments

S1: Overview

Analyze transverse centroid and envelope properties of an unbunched $(\partial/\partial z = 0)$ beam



Transverse averages:

$$\langle \cdots \rangle_\perp \equiv \frac{\int d^2x_\perp \int d^2x'_\perp \; \cdots f_\perp}{\int d^2x_\perp \int d^2x'_\perp \; f_\perp}$$

Centroid:

$$X = \langle x \rangle_\perp$$
$$Y = \langle y \rangle_\perp$$

x- and y-coordinates of beam "center of mass"

Envelope: (edge measure)

$$r_x = 2\sqrt{\langle (x - X)^2 \rangle_\perp}$$
$$r_y = 2\sqrt{\langle (y - Y)^2 \rangle_\perp}$$

x- and y-principal axis radii of an elliptical beam envelope
- Apply to general $f_\perp$ but base on uniform density $f_\perp$
- Factor of 2 results from dimensionality (diff 1D and 3D)

Oscillations in the statistical beam centroid and envelope radii are the *lowest-order* collective responses of the beam

Centroid Oscillations: Associated with errors and are suppressed to the extent possible:
- Error Sources:
  - Beam distribution assymetries (even emerging from injector)
  - Dipole bending terms from applied field optics (due to field error or mech misalignment)
  - Imperfect mechanical alignment
- Exception: When the beam is kicked (insertion or extraction) into or out of a transport channel as is often done in rings

Envelope Oscillations: Can have two components in periodic focusing lattices

1) Matched Envelope: Periodic "flutter" synchronized to period of focusing lattice to yield net focusing
   - Properly tuned flutter essential in Alternating Gradient quadrupole lattices

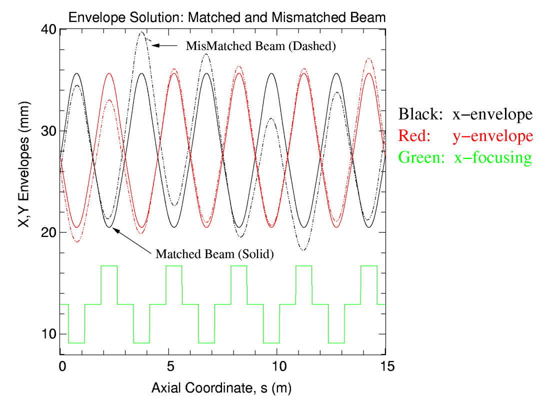
2) Mismatched Envelope: Excursions deviate from matched flutter motion and are seeded/driven by errors

Limiting maximum beam-edge excursions is desired for economical transport
- Reduces cost by Limiting material volume needed to transport an intense beam

Mismatched beams have larger envelope excursions and have more collective stability and beam halo problems since mismatch adds another source of free energy that can drive statistical increases in particle amplitudes

(see: J.J. Barnard lectures on Envelopes and Halo)

Example: FODO Quadrupole Transport Channel



- Larger machine aperture is needed to confine a mismatched beam

Centroid and Envelope oscillations are the *most important collective modes* of an intense beam

- Force balances based on matched beam envelope equation predict scaling of transportable beam parameters
  - Used to design transport lattices
- Instabilities in beam centroid and/or envelope oscillations can prevent reliable transport
  - Parameter locations of instability regions should be understood and avoided in machine design/operation

Although it is *necessary* to avoid envelope and centroid instabilities in designs, it is not alone *sufficient* for effective machine operation

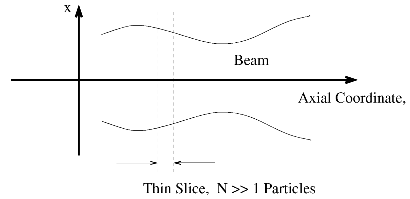
- Higher-order kinetic and fluid instabilities not expressed in the low-order envelope models can degrade beam quality and control and must also be evaluated
  - To be covered (see: S.M. Lund, lectures on Kinetic Stability)

S2: Derivation of Transverse Centroid and Envelope Equations of Motion

Analyze centroid and envelope properties of an unbunched $(\partial/\partial z = 0)$ beam

Transverse Statistical Averages:

Let N be the number of particles in a thin axial slice of the beam at axial coordinate s.



Averages can be equivalently defined in terms of the discreet particles making up the beam or the continuous model transverse Vlasov distribution function:

particles: $\langle \cdots \rangle_\perp \equiv \frac{1}{N} \sum_{i=1}^{N} \Big|_{\text{slice}} \cdots$

distribution: $\langle \cdots \rangle_\perp \equiv \frac{\int d^2x_\perp \int d^2x'_\perp \cdots f_\perp}{\int d^2x_\perp \int d^2x'_\perp \, f_\perp}$

- Averages can be generalized to include axial momentum spread

Transverse Particle Equations of Motion

Consistent with earlier analysis [lectures on Transverse Particle Dynamics], take:

$$x'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)} x' + \kappa_x x = -\frac{q}{m\gamma_b^3\beta_b^2c^2}\frac{\partial\phi}{\partial x}$$

$$y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)} y' + \kappa_y y = -\frac{q}{m\gamma_b^3\beta_b^2c^2}\frac{\partial\phi}{\partial y}$$

$$\nabla_\perp^2\phi = \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)\phi = -\frac{\rho}{\epsilon_0}$$

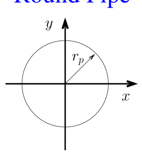
$$\rho = q\int d^2x'_\perp \, f_\perp \qquad \phi|_{\text{aperture}} = 0$$

Assume:

- Unbunched beam
- No axial momentum spread
- Linear applied focusing fields described by κ_x, κ_y
- Possible acceleration, $\gamma_b\beta_b$ need not be constant

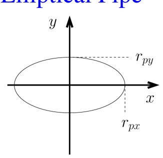
Various apertures are possible influence solution for ϕ. Some simple examples:

Round Pipe



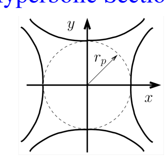
Linac magnetic quadrupoles, acceleration cells, ...

Elliptical Pipe



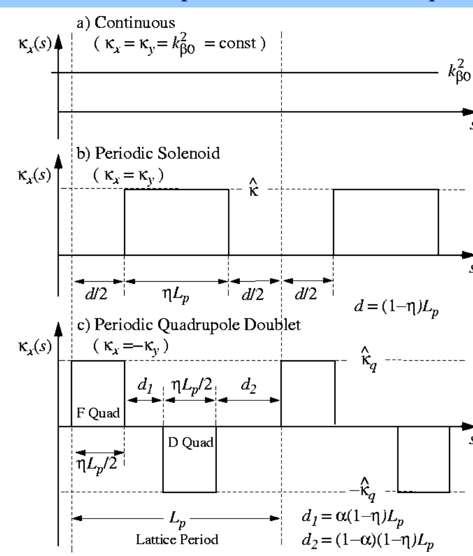
In rings with dispersion: in drifts, magnetic optics, ...

Hyperbolic Sections



Electric quadrupoles

Review: Focusing lattices we will take in examples: Continuous and piecewise constant periodic solenoid and quadrupole doublet



Lattice Period L_p

Occupancy η

$\eta \in [0, 1]$

Solenoid description carried out implicitly in Larmor frame [see: S.M. Lund lectures on Transverse Particle Dynamics]

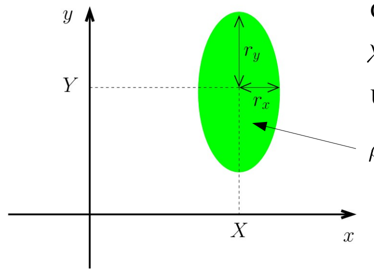
Syncopation Factor α

$\alpha \in [0, \frac{1}{2}]$

$\alpha = \frac{1}{2} \implies FODO$

Distribution Assumptions

To lowest order, linearly focused intense beams are expected to be nearly uniform in density within the core of the beam out to an edge where the density falls rapidly to zero



Charge conservation requires:

$$\lambda = \text{const}$$

Uniform density in beam:

$$\rho = \frac{\lambda}{\pi r_x r_y}$$

$$\rho(x,y) = q\int d^2x'_\perp \; f_\perp \simeq \begin{cases} \frac{\lambda}{\pi r_x r_y}, & (x-X)^2/r_x^2 + (y-Y)^2/r_y^2 < 1 \\ 0, & (x-X)^2/r_x^2 + (y-Y)^2/r_y^2 > 1 \end{cases}$$

$$\lambda = q\int d^2x_\perp \int d^2x'_\perp \; f_\perp \; = \int d^2x \; \rho \; = \text{const}$$

Comments:

- Nearly uniform density out to a sharp beam edge expected for near equilibrium structure beam with strong space-charge due to Debye screening
  - see: S.M. Lund, lectures on Transverse Equilibrium Distributions
- Simulations support that uniform density model is a good approximation for stable non-equilibrium beams when space-charge is high
- Assumption of a fixed form of distribution essentially closes the infinite hierarchy of moments that are needed to describe a general beam distribution
  - Need only describe shape/edge and center for uniform density beam to fully specify the distribution!
  - Analogous to closures of fluid theories using assumed equations of state etc.

Self-Field Calculation

Temporarily, we will consider an *arbitrary* beam charge distribution within an arbitrary aperture to formulate the problem.

Electrostatic field of a line charge in free-space

$$\mathbf{E}_\perp = \frac{\lambda_0}{2\pi\epsilon_0}\frac{(\mathbf{x}_\perp - \tilde{\mathbf{x}})}{|\mathbf{x}_\perp - \tilde{\mathbf{x}}|^2}$$

$\lambda_0 =$ line charge

$\mathbf{x}_\perp = \tilde{\mathbf{x}} =$ coordinate of charge

Resolve the field of the beam into direct (free space) and image terms:

$$\mathbf{E}^s_\perp = -\frac{\partial\phi}{\partial\mathbf{x}_\perp} = \mathbf{E}^d_\perp + \mathbf{E}^i_\perp$$

and superimpose free-space solutions for direct and image contributions

Direct Field

$$\mathbf{E}^d_\perp(\mathbf{x}_\perp) = \frac{1}{2\pi\epsilon_0}\int d^2\tilde{x}_\perp \frac{\rho(\tilde{\mathbf{x}}_\perp)(\mathbf{x}_\perp - \tilde{\mathbf{x}}_\perp)}{|\mathbf{x}_\perp - \tilde{\mathbf{x}}_\perp|^2}$$

$\rho(\mathbf{x}) =$ beam charge density

Image Field

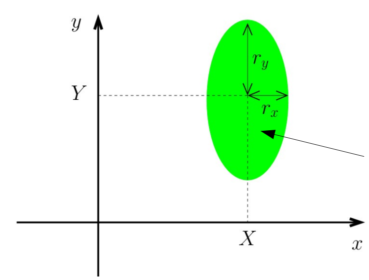
$$\mathbf{E}^i_\perp(\mathbf{x}_\perp) = \frac{1}{2\pi\epsilon_0}\int d^2\tilde{x}_\perp \frac{\rho^i(\tilde{\mathbf{x}}_\perp)(\mathbf{x}_\perp - \tilde{\mathbf{x}}_\perp)}{|\mathbf{x}_\perp - \tilde{\mathbf{x}}_\perp|^2}$$

$\rho^i(\mathbf{x}) =$ beam image charge density induced on aperture

Direct Field:

The direct field solution for a uniform density beam in free-space was calculated for the KV equilibrium distribution

- see: S.M. Lund, lectures on Transverse Equilibrium Distributions



Uniform density in beam:

$$\rho = \frac{\lambda}{\pi r_x r_y} = \text{const}$$

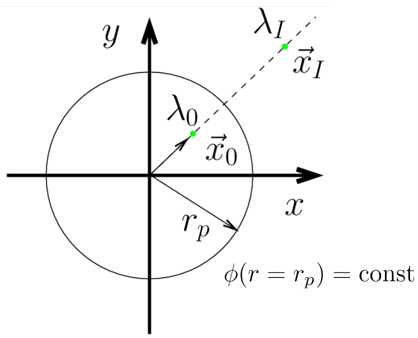
$$E^d_x = \frac{\lambda}{\pi\epsilon_0}\frac{x-X}{(r_x+r_y)r_x}$$

$$E^d_y = \frac{\lambda}{\pi\epsilon_0}\frac{y-Y}{(r_x+r_y)r_y}$$

Expressions are valid only within the elliptical density beam -- where they will be applied in taking averages

Image Field:

Image structure depends on the aperture. Assume a round pipe (most common case) for simplicity.



$\lambda_I = -\lambda_0$ image charge

$\mathbf{x}_I = \frac{r_p^2}{|\mathbf{x}_0|^2}\mathbf{x}_0$ image location

Will be derived in the the problem sets.

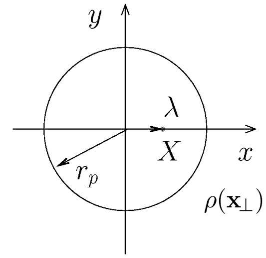
superimpose all images of beam:

$$\mathbf{E}^i_\perp(\mathbf{x}_\perp) = -\frac{1}{2\pi\epsilon_0}\int_{\text{pipe}} d^2\tilde{x}_\perp \frac{\rho(\tilde{\mathbf{x}}_\perp)(\mathbf{x}_\perp - r_p^2\tilde{\mathbf{x}}_\perp/|\tilde{\mathbf{x}}_\perp|^2)}{|\mathbf{x}_\perp - r_p^2\tilde{\mathbf{x}}_\perp/|\tilde{\mathbf{x}}_\perp|^2|^2}$$

- Difficult to calculate even for ρ corresponding to a uniform density beam

Examine limits of the image field to build intuition on the range of properties:

1) Line charge along *x*-axis:



choose coordinates to make true

$$\rho(\mathbf{x}_\perp) = \lambda\delta(\mathbf{x}_\perp - X\hat{\mathbf{x}})$$

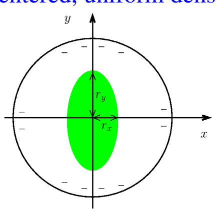
Plug this density in the image charge expression for a round-pipe aperture:

- Need only evaluate at $\mathbf{x}_\perp = X\hat{\mathbf{x}}$ since beam is at that location

$$\mathbf{E}^i_\perp(\mathbf{x}_\perp = X\hat{\mathbf{x}}) = \frac{\lambda}{2\pi\epsilon_0(r_p^2/X - X)}\hat{\mathbf{x}}$$

- Generates nonlinear field at position of direct charge
- Field creates attractive force between direct and image charge

2) Centered, uniform density elliptical beam:



$$\rho(\mathbf{x}_\perp) = \frac{\lambda}{\pi r_x r_y}\begin{cases}\frac{\lambda}{\pi r_x r_y}, & x^2/r_x^2 + y^2/r_y^2 < 1\\ 0, & x^2/r_x^2 + y^2/r_y^2 > 1\end{cases}$$

Expand using complex coordinates starting from the general image expression:

- Image field is in vacuum aperture so complex methods help calculation

$$\underline{E}^{i*} = E^i_x - iE^i_y = \sum_{n=2,4,\cdots}^{\infty} \underline{c}_n \underline{z}^{n-1} \qquad \underline{c}_n = \frac{1}{2\pi\epsilon_0}\int_{\text{pipe}} d^2x_\perp\ \rho(\mathbf{x}_\perp)\frac{(x-iy)^n}{r_p^{2n}}$$

$$\underline{z} = x + iy \qquad = \frac{\lambda n!}{2\pi\epsilon_0 2^n(n/2+1)!(n/2)!}\left(\frac{r_x^2 - r_y^2}{r_p^4}\right)^{n/2}$$

$$i = \sqrt{-1}$$

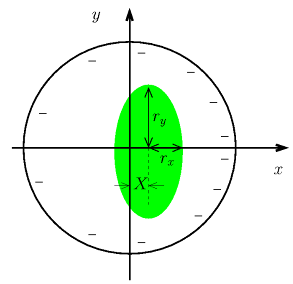
The linear (*n* = 2) components of this expansion give:

$$E^i_x = \frac{\lambda}{8\pi\epsilon_0}\frac{r_x^2 - r_y^2}{r_p^4}x, \qquad E^i_y = -\frac{\lambda}{8\pi\epsilon_0}\frac{r_x^2 - r_y^2}{r_p^4}y$$

- Rapidly vanish (higher order terms more rapid) as beam becomes more round
- Case will be analyzed further in the problem sets

3) Uniform density elliptical beam with a small displacement along the *x*-axis:

$Y = 0 \qquad |X|/r_p \ll 1$



Expand using complex coordinates starting from the general image expression:

- Use complex coordinates to simplify calculation
  E.P. Lee, E. Close, and L. Smith, Nuclear Instruments and Methods, 1126 (1987)
- Expressions become even more complicated with simultaneous *x*- and *y*-displacements and more complicated aperture geometries

Leading order terms expanded in $|X|/r_p$ without assuming small ellipticity obtain:

$$E_x^i = \frac{\lambda}{2\pi\epsilon_0 r_p^2}[f\cdot(x-X)+g\cdot X] + \Theta\left(\frac{X}{r_p}\right)^3$$

$$E_y^i = -\frac{\lambda}{2\pi\epsilon_0 r_p^2} f\cdot y + \Theta\left(\frac{X}{r_p}\right)^3$$

Where f and g are focusing and bending coefficients that can be calculated in terms of $X,\ Y,\ \ r_x,\ r_y$ (which all may vary in s) as:

Focusing Term:

$$f = \frac{r_x^2-r_y^2}{4r_p^2} + \frac{X^2}{r_p^2}\left[1+\frac{3}{2}\left(\frac{r_x^2-r_y^2}{r_p^2}\right)+\frac{3}{8}\left(\frac{r_x^2-r_y^2}{r_p^2}\right)^2\right]$$

Bending Term:

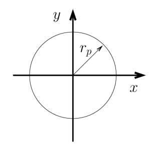
$$g = 1 + \frac{r_x^2-r_y^2}{4r_p^2} + \frac{X^2}{r_p^2}\left[1+\frac{3}{4}\left(\frac{r_x^2-r_y^2}{r_p^2}\right)+\frac{1}{8}\left(\frac{r_x^2-r_y^2}{r_p^2}\right)^2\right]$$

- Expressions become even more complicated with simultaneous x- and y-displacements and more complicated aperture geometries

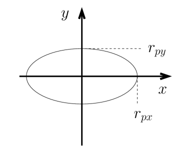
Comments on images:

- Sign is generally such that it will tend to increase beam displacements
  - Also (usually) weak linear focusing corrections for an elliptical beam
- Can be very difficult to calculate explicitly
  - Even for simple case of circular pipe
  - Special cases of simple geometry formulas can give idea on scaling
  - Generally suppress just by making the beam small relative to characteristic aperture dimensions and keeping the beam steered near-axis
  - Simulations typically applied
- Depend strongly on the aperture geometry
  - Generally varies as a function of s in the machine aperture due to changes in accelerator lattice elements and/or as beam symmetries evolve

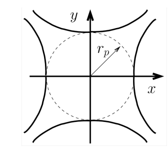
Round Pipe



Elliptical Pipe



Hyperbolic Sections



Coupled centroid and envelope equations of motion

Consistent with the assumed structure of the distribution (uniform density elliptical beam), denote:

Beam Centroid:

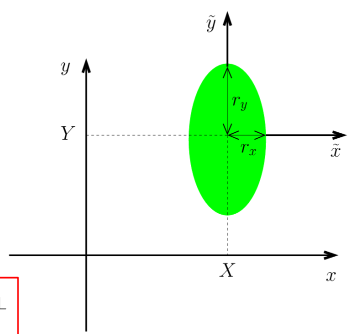
$$X \equiv \langle x\rangle_\perp \qquad X' = \langle x'\rangle_\perp$$
$$Y \equiv \langle y\rangle_\perp \qquad Y' = \langle y'\rangle_\perp$$

Coordinates with respect to centroid:

$$\tilde{x} \equiv x - X \qquad \tilde{x}' = x' - X'$$
$$\tilde{y} \equiv y - Y \qquad \tilde{y}' = y' - Y'$$

Envelope Edge Radii:

$$r_x \equiv 2\sqrt{\langle\tilde{x}^2\rangle_\perp} \qquad r_x' = 2\langle\tilde{x}\tilde{x}'\rangle_\perp/\langle\tilde{x}^2\rangle_\perp$$
$$r_y \equiv 2\sqrt{\langle\tilde{y}^2\rangle_\perp} \qquad r_y' = 2\langle\tilde{y}\tilde{y}'\rangle_\perp/\langle\tilde{y}^2\rangle_\perp$$



With the *assumed* uniform elliptical beam, all moments can be calculated in terms of: $X,\ Y \quad r_x,\ r_y$

- Such truncations follow whenever the form of the distribution is "frozen"

Derive centroid equations: First use the self-field resolution for a uniform density beam, then the equations of motion for a particle within the beam are:

$$x'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}x' + \kappa_x x - \frac{2Q}{(r_x+r_y)r_x}(x-X) = \frac{q}{m\gamma_b^3\beta_b^2c^2}E_x^i$$

$$y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}y' + \kappa_y y - \frac{2Q}{(r_x+r_y)r_y}(y-Y) = \frac{q}{m\gamma_b^3\beta_b^2c^2}E_y^i$$

Direct Terms Image Terms

Perveance:

$$Q \equiv \frac{q\lambda}{2\pi\epsilon_0 m\gamma_b^3\beta_b^2c^2}$$

(not necessarily constant if beam accelerates)

average equations using: $\langle x'\rangle_\perp = \langle x\rangle_\perp' = X'$ etc., to obtain:

Centroid Equations:

$$X'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}X' + \kappa_x X = Q\left[\frac{2\pi\epsilon_0}{\lambda}\langle E_x^i\rangle_\perp\right]$$

$$Y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}Y' + \kappa_y Y = Q\left[\frac{2\pi\epsilon_0}{\lambda}\langle E_y^i\rangle_\perp\right]$$

Note: the electric image field will cancel the coefficient $2\pi\epsilon_0/\lambda$

- $\langle E_x^i\rangle_\perp$ will generally depend on: $X,\ \ Y$ and $r_x,\ \ r_y$

To derive equations of motion for the envelope radii, first subtract the centroid equations from the particle equations of motion ($\tilde{x} \equiv x - X$) to obtain:

$$\tilde{x}'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}\tilde{x}' + \kappa_x\tilde{x} - \frac{2Q\tilde{x}}{(r_x+r_y)r_x} = \frac{q}{m\gamma_b^3\beta_b^2c^2}\left[E_x^i - \langle E_x^i\rangle_\perp\right]$$
$$\tilde{y}'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}\tilde{y}' + \kappa_y\tilde{y} - \frac{2Q\tilde{y}}{(r_x+r_y)r_x} = \frac{q}{m\gamma_b^3\beta_b^2c^2}\left[E_y^i - \langle E_y^i\rangle_\perp\right]$$

Differentiate the equation for the envelope radius (y-equations analogous):

$$r_x = 2\langle\tilde{x}^2\rangle_\perp^{1/2} \implies r_x' = \frac{2\langle\tilde{x}\tilde{x}'\rangle_\perp}{\langle\tilde{x}^2\rangle_\perp^{1/2}} = \frac{4\langle\tilde{x}\tilde{x}'\rangle_\perp}{r_x}$$

Define (motivated the KV equilibrium results) a statistical rms edge emittance:

$$\varepsilon_x \equiv 4\varepsilon_{x,\text{rms}} \equiv 4\left[\langle\tilde{x}^2\rangle_\perp\langle\tilde{x}'^2\rangle_\perp - \langle\tilde{x}\tilde{x}'\rangle_\perp^2\right]^{1/2}$$

Differentiate the equation for r_x' again and use the emittance definition:

$$r_x'' = 4\frac{\langle\tilde{x}\tilde{x}''\rangle_\perp}{r_x} + \frac{16[\langle\tilde{x}^2\rangle_\perp\langle\tilde{x}'^2\rangle_\perp - \langle\tilde{x}\tilde{x}'\rangle_\perp^2]}{r_x^3}$$
$$= 4\frac{\langle\tilde{x}\tilde{x}''\rangle_\perp}{r_x} + \frac{\varepsilon_x^2}{r_x^3}$$

and then employ the equations of motion to eliminate $\tilde{x}''$ in $\langle\tilde{x}\tilde{x}''\rangle_\perp$ to obtain:

Envelope Equations:

$$r_x'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}r_x' + \kappa_x r_x - \frac{2Q}{r_x+r_y} - \frac{\varepsilon_x^2}{r_x^3} = 8Q\left[\frac{\pi\epsilon_0}{\lambda}\langle\tilde{x}E_x^i\rangle_\perp\right]$$
$$r_y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}r_x' + \kappa_y r_y - \frac{2Q}{r_x+r_y} - \frac{\varepsilon_y^2}{r_y^3} = 8Q\left[\frac{\pi\epsilon_0}{\lambda}\langle\tilde{y}E_y^i\rangle_\perp\right]$$

- $\langle\tilde{x}E_x^i\rangle_\perp$ will generally depend on: $X,\ \ Y$ and $r_x,\ \ r_y$

Comments on Centroid/Envelope equations:

- Centroid and envelope equations are *coupled* and must be solved simultaneously when image terms on the RHS cannot be neglected
- Image terms contain nonlinear terms that can be difficult to evaluate explicitly
  - Aperture geometry changes image correction
- The formulation is not self-consistent because a frozen form (uniform density) charge profile is assumed
  - Uniform density choice motivated by KV results and Debye screening see: S.M. Lund, lectures on Transverse Equilibrium Distributions
  - The assumed distribution form not evolving represents a fluid model closure
  - Generally find with simulations that uniform density frozen form distribution models can provide reasonably accurate approximate models for centroid and envelope evolution

Comments on Centroid/Envelope equations (Continued):

- Constant (normalized when accelerating) emittances are generally assumed
  - For strong space charge emittance terms small and limited emittance evolution does not strongly influence evolution outside of final focus
  - See: S.M. Lund, lectures on Transverse Particle Dynamics and Transverse Kinetic Theory to motivate when this works well

$\beta_b,\ \gamma_b,\ \lambda$ s-variation set by acceleration schedule

$$\varepsilon_{nx} = \gamma_b\beta_b\varepsilon_x = \text{const}$$
$$\varepsilon_{ny} = \gamma_b\beta_b\varepsilon_y = \text{const}$$
$\longrightarrow$ used to calculate $\varepsilon_x,\ \varepsilon_y$

$$Q = \frac{q\lambda}{2\pi m\epsilon_0\gamma_b^3\beta_b^2c^2}$$

S3: Centroid Equations of Motion
Single Particle Limit: Oscillation and Stability Properties

Neglect image charge terms, then the centroid equation of motion becomes:

$$X'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}X' + \kappa_x X = 0$$
$$Y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}Y' + \kappa_y Y = 0$$

- Usual Hill's equation with acceleration term
- Single particle form. Apply results from S.M. Lund lectures on Transverse Particle Dynamics: phase amplitude methods, Courant-Snyder invariants, and stability bounds, ...

Assume that applied lattice focusing is tuned for constant phase advances with normalized coordinates and/or that acceleration is weak and can be neglected with Then single particle stability results give immediately:

$$\frac{1}{2}|\text{Tr}\ \mathbf{M}_x(s_i+L_p|s_i)| \leq 1$$
$$\frac{1}{2}|\text{Tr}\ \mathbf{M}_y(s_i+L_p|s_i)| \leq 1$$
$$\iff \quad \sigma_{0x} < 180^\circ, \quad \sigma_{0y} < 180^\circ$$

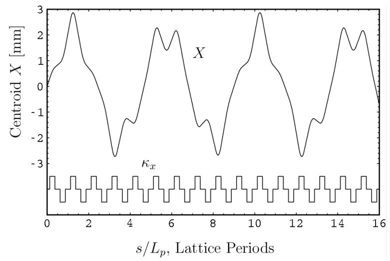
centroid stability
1st stability condition

/// Example: FODO channel centroid evolution

Mid-drift launch:

$X(0) = 1$ mm

$X'(0) = 1$ mrad



lattice/beam parameters:

$\beta_b = \text{const}$

$\sigma_{0x} = 80°$

$L_p = 0.5$ m

$\eta = 0.5$

- Centroid exhibits expected characteristic stable betatron oscillations
- Motion in *y*-plane analogous

///

Effect of Driving Errors

The reference orbit is ideally tuned for zero centroid excursions. But there will *always* be driving errors that can cause the centroid oscillations to accumulate with beam propagation distance:

$$X'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}X' + \sum_n \frac{G_n}{G_0}\kappa_n(s)X = \sum_n \frac{G_n}{G_0}\kappa_n(s)\Delta_{xn}$$

$\kappa_q(s) = \sum_n \kappa_n(s)$ $\kappa_n(s)$ nominal gradient function, *n*th quadrupole

$\frac{G_n}{G_0} =$ *n*th quadrupole gradient error (unity for no error; *s*-varying)

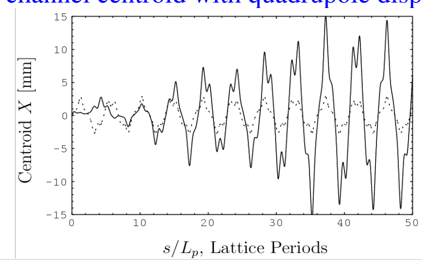
$\Delta_{xn} =$ *n*th quadrupole transverse displacement error (*s*-varying)

/// Example: FODO channel centroid with quadrupole displacement errors

$\frac{G_n}{G_0} = 1$

$\Delta_{xn} = [-0.5, 0.5]$ mm (uniform dist)

same lattice as previous



solid – with errors
dashed – no errors

///

Errors will result in a characteristic random walk increase in oscillation amplitude due to the (generally random) driving terms.

Control by:

- Synthesize small applied dipole fields to regularly steer the centroid back on-axis to the reference trajectory: X = 0 = Y, X' = 0 = Y'
- Fabricate and align focusing elements with higher precision
- Employ a sufficiently large aperture to contain the oscillations and limit detrimental nonlinear image charge effects

Economics dictates the optimal strategy

- Usually sufficient control achieved by a combination of methods

Effects of Image Charges

Model the beam as a displaced line-charge in a circular aperture. Then using the previously derived image charge field, the equations of motion reduce to:

$$X'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}X' + \kappa_x X = \frac{QX}{r_p^2 - X^2}$$

examine oscillation along *x*-axis

$$\frac{QX}{r_p^2 - X^2} \simeq \frac{Q}{r_p^2}X + \frac{Q}{r_p^4}X^3$$

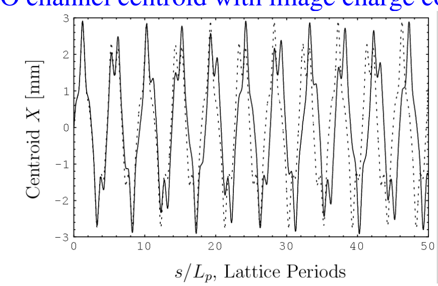
linear correction — Nonlinear correction (smaller)

Example: FODO channel centroid with image charge corrections

$r_p = 30$ mm

$Q = 2 \times 10^{-4}$

same lattice as previous



solid – with images
dashed – no images

Main effect of images is generally an accumulated phase error of the centroid orbit since, generally the centroid error oscillations are not "matched" orbits and errors are not regularly "undone"

- This will complicate extrapolations of errors over many lattice periods

Control by:

- Keeping centroid displacements X, Y small by correcting
- Make aperture (pipe radius) larger

General Comments:

- Images contributions to centroid excursions generally less problematic than misalignment errors in focusing elements
- More detailed analysis show that the coupling of the envelope radii $r_x,\ r_y$ to the centroid evolution in X, Y is often weak
- Fringe fields are more important for accurate calculation of centroid orbits since orbits are not part of a matched lattice
  - Non-ideal orbits are poorly tuned to lattice and become more sensitive to the precise phase of impulses
- Over long path lengths many nonlinear terms can influence results
- Lattice errors are not often known so one must often analyze characteristic error distributions to see if centroids measured are consistent with expectations

S4: Envelope Equations of Motion

Overview: Reduce equations of motion for $r_x,\ r_y$

- Generally found that couplings to centroid coordinates $X\ \ Y$ are weak
  - Centroid ideally zero in a well tuned system
- Envelope eqns are most important in designing transverse focusing systems
  - Expresses average radial force balance (see following discussion)
  - Can be difficult to analyze analytically for scaling properties
  - "Systems" codes generally written using envelope equations, stability criteria, and practical engineering constraints
- Instabilities of the envelope equations in periodic focusing lattices must be avoided in machine operation
  - Instabilities are strong and real: not washed out with realistic distributions without frozen form
  - Represent lowest order "KV" modes of a full kinetic theory
- Previous derivation of envelope equations relied on Courant-Snyder invariants in linear applied and self-fields. Analysis shows that the same force balances result for a uniform elliptical beam with no image couplings.
  - Debye screening arguments suggest assumed uniform density model taken should be a good approximation for intense space-charge

KV/rms Envelope Equations: Properties of Terms

The envelope equation reflects low-order force balances:

$$r_x'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}r_x' + \kappa_x r_x - \frac{2Q}{r_x+r_y} - \frac{\varepsilon_x^2}{r_x^3} = 0$$

$$r_y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}r_y' + \kappa_y r_y - \frac{2Q}{r_x+r_y} - \frac{\varepsilon_y^2}{r_y^3} = 0$$

Terms:	Applied Acceleration	Applied Focusing	Space-Charge Defocusing	Thermal Defocusing
	Lattice	Lattice	Perveance	Emittance

The "acceleration schedule" specifies both $\gamma_b\beta_b$ and λ then the equations are integrated with:

$$\gamma_b\beta_b\varepsilon_x = \text{const} \qquad \gamma_b\beta_b\varepsilon_y = \text{const}$$

normalized emittance conservation

$$Q = \frac{q\lambda}{2\pi\epsilon_0 m\gamma_b^3\beta_b^2c^2}$$

specified perveance

Reminder: It was shown for a coasting beam that the envelope equations remain valid for elliptic charge densities suggesting more general validity

[Sacherer, IEEE Trans. Nucl. Sci. **18**, 1101 (1971), J.J. Barnard, Intro. Lectures]

For any beam with elliptic symmetry charge density in each transverse slice:

$$\rho = \rho\left(\frac{x^2}{r_x^2} + \frac{y^2}{r_y^2}\right)$$

Based on:

$$\langle x\frac{\partial\phi}{\partial x}\rangle_\perp = -\frac{\lambda}{4\pi\epsilon_0}\frac{r_x}{r_x+r_y}$$

see J.J. Barnard, Intro. Lectures

the KV envelope equations

$$r_x''(s) + \kappa_x(s)r_x(s) - \frac{2Q}{r_x(s)+r_y(s)} - \frac{\varepsilon_x^2(s)}{r_x^3(s)} = 0$$

$$r_y''(s) + \kappa_y(s)r_y(s) - \frac{2Q}{r_x(s)+r_y(s)} - \frac{\varepsilon_y^2(s)}{r_y^3(s)} = 0$$

remain valid when (averages taken with the full distribution):

$$Q = \frac{q\lambda}{2\pi\epsilon_0 m\gamma_b^3\beta_b^2c^2} = \text{const} \qquad \lambda = q\int d^2x_\perp\ \rho\ = \ \text{const}$$

$$r_x = 2\langle x^2\rangle_\perp^{1/2} \qquad \varepsilon_x = 4[\langle x^2\rangle_\perp\langle x'^2\rangle_\perp - \langle xx'\rangle_\perp^2]^{1/2}$$

$$r_y = 2\langle y^2\rangle_\perp^{1/2} \qquad \varepsilon_y = 4[\langle y^2\rangle_\perp\langle y'^2\rangle_\perp - \langle yy'\rangle_\perp^2]^{1/2}$$

- Evolution changes often small in $\varepsilon_x,\ \varepsilon_y$

Properties of Envelope Equation Terms:

Applied Focusing: $\kappa_x r_x, \quad \kappa_y r_y$ and Acceleration: $\frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)} r_x', \quad \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)} r_y'$

- Analogous to single particle orbit terms
- Contributions to beam envelope essentially the same as in single particle case
- Have strong *s* dependence, *can be both focusing and defocusing*
  - Act only in focusing elements and acceleration gaps

Perveance: $\frac{2Q}{r_x + r_y}$

- Acts continuously in *s*, *always defocusing*
- Becomes stronger (relatively to other terms) when the beam expands in cross-sectional area

Emittance: $\frac{\varepsilon_x^2}{r_x^3}$

- Acts continuously in *s*, *always defocusing*
- Becomes stronger (relatively to other terms) when the beam becomes small in cross-sectional area
- Scaling makes clear why it is necessary to inhibit emittance growth for applications where small spots are desired on target

As the beam expands, perveance term will eventually dominate emittance term:
[see: Lund and Bukh, PRSTAB **7**, 024801 (2004)]

Free expansion $(\kappa_x = \kappa_y = 0)$

Initial conditions:

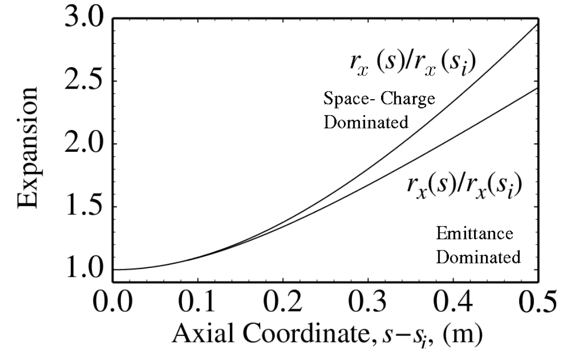
$$r_x(s_i) = r_y(s_i) \qquad \frac{Q}{r_x(s_i)} = \frac{\varepsilon_x^2}{r_x^3(s_i)}$$

$$r_x'(s_i) = r_y'(s_i) = 0 \qquad Q = \frac{\varepsilon_x^2}{r_x^2(s_i)} = 10^{-3}$$

Cases:

Space-Charge Dominated: $\varepsilon_x = 0$

Emittance Dominated: $Q = 0$



See next page: solution is analytical in bounding limits shown

Parameters are chosen such that initial defocusing forces in two limits are equal to compare case

For an emittance dominated beam in free-space, the envelope equation becomes:

$$\frac{Q}{r_x + r_y} \ll \frac{\varepsilon_{x,y}^2}{r_{x,y}^3} \quad \Longrightarrow \quad r_j'' - \frac{\varepsilon_j^2}{r_j^3} = 0 \qquad j = x, y$$

The envelope Hamiltonian gives:

$$\frac{1}{2} r_j'^2 + \frac{\varepsilon_j^2}{2 r_j^2} = \text{const}$$

which can be integrated from the initial envelope at $s = s_i$ to show that:

Emittance Dominated Free-Expansion $(Q = 0)$

$$r_j(s) = r_j(s_i)\sqrt{1 + \frac{2r_j'(s_i)}{r_j(s_i)}(s - s_i) + \left[1 + \frac{r_j^2(s_i) r_j'^2(s_i)}{\varepsilon_j^2}\right] \frac{\varepsilon_j^2}{r_j^4(s_i)}(s - s_i)^2}$$

$j = x, y$

Conversely, for a space-charge dominated beam in free-space, the envelope equation becomes:

$$\frac{Q}{r_x + r_y} \gg \frac{\varepsilon_{x,y}^2}{r_{x,y}^3} \quad \Longrightarrow \quad \begin{aligned} & r_+'' - \frac{Q}{r_+} = 0 \\ & r_-'' = 0 \end{aligned} \qquad r_\pm \equiv \frac{1}{2}(r_x \pm r_y)$$

The equations of motion

$$r_+'' - \frac{Q}{r_+} = 0$$

$$r_-'' = 0$$

can be integrated from the initial envelope at $s = s_i$ to show that:

- r_- equation solution trivial
- r_+ equation solution exploits Hamiltonian $\frac{1}{2} r_+'^2 - Q \ln r_+ = \text{const}$

Space-Charge Dominated Free-Expansion $(\varepsilon_x = \varepsilon_y = 0)$

$$r_+(s) = r_+(s_i) \exp\left(-\frac{r_+'^2(s_i)}{2Q} + \left[\text{erfi}^{-1}\left\{\text{erfi}\left[\frac{r_+'(s_i)}{\sqrt{2Q}}\right] + \sqrt{\frac{2Q}{\pi}} e^{\frac{r_+'^2(s_i)}{2Q}} \frac{(s - s_i)}{r_+(s_i)}\right\}\right]^2\right)$$

$$r_-(s) = r_-(s_i) + r_-'(s_i)(s - s_i)$$

$$r_\pm = \frac{1}{2}(r_x \pm r_y)$$

Imaginary Error Function

$$\text{erfi}(z) \equiv \frac{\text{erf}(iz)}{i} \equiv \frac{2}{\sqrt{\pi}} \int_0^z dt\ \exp(t^2)$$

$$i \equiv \sqrt{-1}$$

The free-space expansion solutions for emittance and space-charge dominated beams will be explored more in the problems

S5: Matched Envelope Solution:

Lund and Bukh, PRSTAB **7**, 024801 (2004)

Neglect acceleration $(\gamma_b \beta_b = \text{const})$ or use transformed variables:

$$r_x''(s) + \kappa_x(s) r_x(s) - \frac{2Q}{r_x(s)+r_y(s)} - \frac{\varepsilon_x^2}{r_x^3(s)} = 0$$

$$r_y''(s) + \kappa_y(s) r_y(s) - \frac{2Q}{r_x(s)+r_y(s)} - \frac{\varepsilon_y^2}{r_y^3(s)} = 0$$

$$r_x(s+L_p) = r_x(s) \qquad r_x(s) > 0$$

$$r_y(s+L_p) = r_y(s) \qquad r_y(s) > 0$$

Matching involves finding specific initial conditions for the envelope to have the periodicity of the lattice:

Find Values of:

$$r_x(s_i) \quad r_x'(s_i)$$
$$r_y(s_i) \quad r_y'(s_i)$$

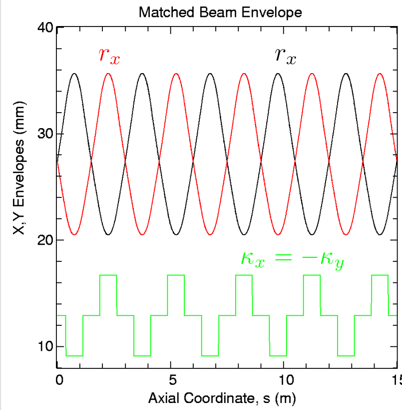
Such That: (periodic)

$$r_x(s_i+L_p) = r_x(s_i) \qquad r_x'(s_i+L_p) = r_x'(s_i)$$
$$r_y(s_i+L_p) = r_y(s_i) \qquad r_y'(s_i+L_p) = r_y'(s_i)$$

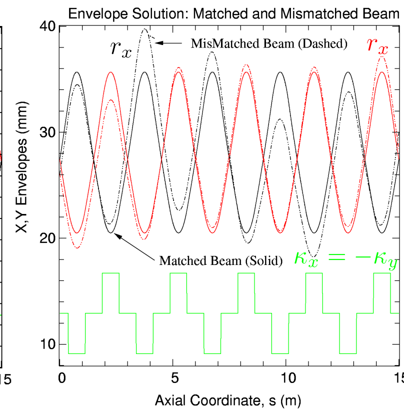
- Typically constructed with numerical root finding from estimated/guessed values
  - Can be surprisingly difficult for complicated lattices and/or strong space-charge
- Iterative technique developed to numerically calculate without root finding
  [Lund, Chilton and Lee, PRSTAB **9**, 064201 (2006)]
  Method exploits Courant-Snyder invariants of depressed orbits within the beam.

Typical Matched vs Mismatched solution for FODO channel:

Matched



Mismatched



The matched beam is the most radially compact solution to the envelope equations rendering it highly important for beam transport

- Matching tends to exploit optics most efficiently to maintain confinement

The matched solution to the KV envelope equations reflects the symmetry of the focusing lattice and must in general be calculated numerically

$$r_x(s+L_p) = r_x(s)$$
$$r_y(s+L_p) = r_y(s)$$
$$\varepsilon_x = \varepsilon_y$$

Parameters

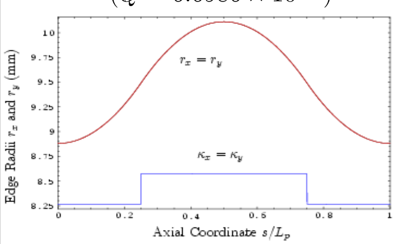
$L_p = 0.5$ m, $\sigma_0 = 80°$, $\eta = 0.5$

$\varepsilon_x = 50$ mm-mrad

$\sigma/\sigma_0 = 0.2$ Perveance Q iterated to obtain matched solution with this tune depression

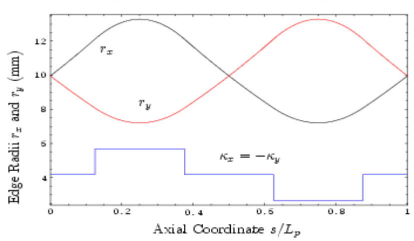
Solenoidal Focusing

$(Q = 6.6986 \times 10^{-4})$



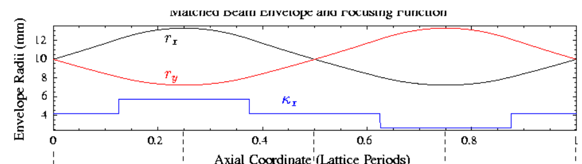
FODO Quadrupole Focusing

$(Q = 6.5614 \times 10^{-4})$

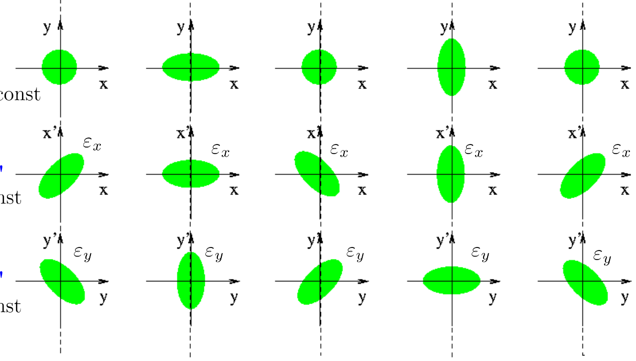


Symmetries of a matched beam are interpreted in terms of a local rms equivalent KV beam and moments/projections of the KV distribution

[see: S.M. Lund, lectures on Transverse Equilibrium Distributions]



Projection



x-y
area: $\pi r_x r_y \neq$ const

x-x'
area: $\pi \varepsilon_x =$ const
(CS Invariant)

y-y'
area: $\pi \varepsilon_y =$ const
(CS Invariant)

S6: Envelope Perturbations:

Lund and Bukh, PRSTAB **7**, 024801 (2004)

In the envelope equations take:

Envelope Perturbations:

$$r_x(s) = r_{xm}(s) + \delta r_x(s)$$
$$r_y(s) = r_{ym}(s) + \delta r_y(s)$$

Matched Envelope (r_{xm}, r_{ym}); Mismatch Perturbations ($\delta r_x, \delta r_y$)

Driving Perturbations:

$$\kappa_x(s) \to \kappa_x(s) + \delta\kappa_x(s) \qquad \kappa_y(s) \to \kappa_y(s) + \delta\kappa_y(s) \qquad \text{Focus}$$
$$Q \to Q + \delta Q(s) \qquad \text{Perveance}$$
$$\varepsilon_x \to \varepsilon_x + \delta\varepsilon_x(s) \qquad \varepsilon_y \to \varepsilon_y + \delta\varepsilon_y(s) \qquad \text{Emittance}$$

Perturbations in envelope radii are about a matched solution:

$$r_{xm}(s+L_p) = r_{xm}(s) \qquad r_{xm}(s) > 0$$
$$r_{ym}(s+L_p) = r_{ym}(s) \qquad r_{ym}(s) > 0$$

Perturbations in envelope radii are small relative to matched solution and driving terms are consistently ordered:

$$r_{xm}(s) \gg |\delta r_x(s)|$$
$$r_{ym}(s) \gg |\delta r_y(s)|$$

$\longleftarrow$ Amplitudes defined in terms of producing small envelope perturbations

- Driving perturbations and distribution errors generate/pump envelope perturbations
  - Arise from many sources: focusing errors, lost particles, emittance growth, ...

The matched solution satisfies:

- Add subscript *m* to denote matched envelope solution and distinguish from other evolutions

$$r_x \to r_{xm} \qquad r_y \to r_{ym}$$

For matched beam envelope with periodicity of lattice

$$r''_{xm}(s) + \kappa_x(s) r_{xm}(s) - \frac{2Q}{r_{xm}(s)+r_{ym}(s)} - \frac{\varepsilon_x^2}{r_{xm}^3(s)} = 0$$
$$r''_{ym}(s) + \kappa_y(s) r_{ym}(s) - \frac{2Q}{r_{xm}(s)+r_{ym}(s)} - \frac{\varepsilon_y^2}{r_{ym}^3(s)} = 0$$
$$r_{xm}(s+L_p) = r_{xm}(s) \qquad r_{xm}(s) > 0$$
$$r_{ym}(s+L_p) = r_{ym}(s) \qquad r_{ym}(s) > 0$$

Matching is usually cast in terms of finding 4 "initial" envelope phase-space values where the envelope solution satisfies the periodicity constraint for specified focusing, perveance, and emittances:

$$r_{xm}(s_i) \qquad r'_{xm}(s_i)$$
$$r_{ym}(s_i) \qquad r'_{ym}(s_i)$$

Linearized Perturbed Envelope Equations:

- Neglect all terms of order δ^2 and higher: $(\delta r_x)^2, \;\; \delta r_x \delta r_y, \;\; \delta Q \delta r_x, \; \cdots$

$$\delta r''_x + \kappa_x \delta r_x + \frac{2Q}{(r_{xm}+r_{ym})^2}(\delta r_x + \delta r_y) + \frac{3\varepsilon_x^2}{r_{xm}^4}\delta r_x = -r_{xm}\delta\kappa_x + \frac{2}{r_{xm}+r_{ym}}\delta Q + \frac{2\varepsilon_x}{r_{xm}^3}\delta\varepsilon_x$$
$$\delta r''_y + \kappa_y \delta r_y + \frac{2Q}{(r_{xm}+r_{ym})^2}(\delta r_x + \delta r_y) + \frac{3\varepsilon_y^2}{r_{ym}^4}\delta r_y = -r_{ym}\delta\kappa_y + \frac{2}{r_{xm}+r_{ym}}\delta Q + \frac{2\varepsilon_y}{r_{ym}^3}\delta\varepsilon_y$$

Homogeneous Equations:

- Linearized envelope equations with driving terms set to zero

$$\delta r''_x + \kappa_x \delta r_x + \frac{2Q}{(r_{xm}+r_{ym})^2}(\delta r_x + \delta r_y) + \frac{3\varepsilon_x^2}{r_{xm}^4}\delta r_x = 0$$
$$\delta r''_y + \kappa_y \delta r_y + \frac{2Q}{(r_{xm}+r_{ym})^2}(\delta r_x + \delta r_y) + \frac{3\varepsilon_y^2}{r_{ym}^4}\delta r_y = 0$$

Martix Form of the Linearized Perturbed Envelope Equations:

$$\frac{d}{ds}\delta\mathbf{R} + \mathbf{K}\cdot\delta\mathbf{R} = \delta\mathbf{P}$$

$$\delta\mathbf{R} \equiv \begin{pmatrix} \delta r_x \\ \delta r'_x \\ \delta r_y \\ \delta r'_y \end{pmatrix}$$ Coordinate vector

$$\mathbf{K} \equiv \begin{pmatrix} 0 & -1 & 0 & 0 \\ k_{xm} & 0 & k_{0m} & 0 \\ 0 & 0 & 0 & -1 \\ k_{0m} & 0 & k_{ym} & 0 \end{pmatrix}$$

Coefficient matrix — Has periodicity of the lattice period

$$k_{0m} = \frac{2Q}{(r_{xm}+r_{ym})^2}$$
$$k_{jm} = \kappa_j + 3\frac{\varepsilon_j^2}{r_{jm}^4} + k_{0m} \qquad j = x, \; y$$

$$\delta\mathbf{P} \equiv \begin{pmatrix} 0 \\ -\delta\kappa_x r_{xm} + 2\frac{\delta Q}{r_{xm}+r_{ym}} + 2\frac{\varepsilon_x\delta\varepsilon_x}{r_{xm}^3} \\ 0 \\ -\delta\kappa_y r_{ym} + 2\frac{\delta Q}{r_{xm}+r_{ym}} + 2\frac{\varepsilon_y\delta\varepsilon_y}{r_{ym}^3} \end{pmatrix}$$ Driving perturbation vector

Expand solution into homogeneous and particular parts:

$$\delta\mathbf{R} = \delta\mathbf{R}_h + \delta\mathbf{R}_p$$

$\delta\mathbf{R}_h =$ homogeneous solution

$\delta\mathbf{R}_p =$ particular solution

$$\frac{d}{ds}\delta\mathbf{R}_h + \mathbf{K}\cdot\delta\mathbf{R}_h = 0 \qquad \frac{d}{ds}\delta\mathbf{R}_p + \mathbf{K}\cdot\delta\mathbf{R}_p = \delta\mathbf{P}$$

Homogeneous Solution: Normal Modes

- Describes normal mode oscillations
- Original analysis by Struckmeier and Reiser [Part. Accel. **14**, 227 (1984)]

Particular Solution: Driven Modes

- Describes action of driving terms
- Characterize in terms of projections on homogeneous response (on normal modes)

Homogeneous solution expressible as a map:

$$\delta \mathbf{R}(s) = \mathbf{M}_e(s|s_i) \cdot \delta \mathbf{R}(s_i)$$
$$\delta \mathbf{R}(s) = (\delta r_x, \delta r'_x, \delta r_y, \delta r'_y)$$
$$\mathbf{M}_e(s|s_i) = \ 4 \times 4 \text{ transfer map}$$

Now 4x4 system, but analogous to the 2x2 analysis of Hill's equation via transfer matrices: see S.M. Lund lectures on Transverse Particle Dynamics

Eigenvalues and eigenvectors of map through one period characterize normal modes and stability properties:

$$\mathbf{M}_e(s_i + L_p|s_i) \cdot \mathbf{E}_n(s_i) = \lambda_n \mathbf{E}_n(s_i)$$

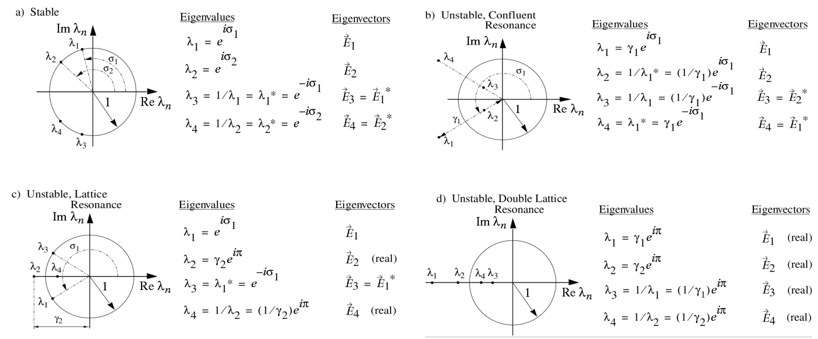
Stability

$$\lambda_n = \gamma_n e^{i\sigma_n} \qquad \begin{array}{l} \sigma_n \rightarrow \text{mode phase advance (real)} \\ \gamma_n \rightarrow \text{mode growth/damp factor (real)} \end{array}$$

Mode Expansion/Launching

$$\delta \mathbf{R}(s_i) = \sum_{n=1}^{4} \alpha_n \mathbf{E}_n(s_i)$$
$$\alpha_n = \text{const (complex)}$$

Eigenvalue/Eigenvector Symmetry Classes:



Symmetry classes of eigenvalues/eigenvectors:

- Determine normal mode symmetries
- Hamiltonian dynamics allow only 4 distinct classes of eigenvalue symmetries
  - See A. Dragt, Lectures on Nonlinear Orbit Dynamics, in Physics of High Energy Particle Accelerators, (AIP Conf. Proc. No. 87, 1982, p. 147)
- Envelope mode symmetries discussed fully in PRSTAB review
- Caution: Textbook by Reiser makes errors in quadrupole mode symmetries and mislabels/identifies dispersion characteristics and branch choices

Pure mode launching conditions:

Launching conditions for distinct normal modes corresponding to the eigenvalue classes illustrated:

$$A_\ell = \text{mode amplitude (real)} \qquad \ell = \ \text{mode index}$$
$$\psi_\ell = \text{mode launch phase (real)} \qquad C.C. = \ \text{complex conjugate}$$

Case	Mode	Launching Condition	Lattice Period Advance
(a) Stable	1 - Stable Osc.	$\delta\mathbf{R}_1 = A_1 e^{i\psi_1}\mathbf{E}_1 + \text{C.C.}$	$\mathbf{M}_e\delta\mathbf{R}_1(\psi_1) = \delta\mathbf{R}_1(\psi_1 + \sigma_1)$
	2 - Stable Osc.	$\delta\mathbf{R}_2 = A_2 e^{i\psi_2}\mathbf{E}_2 + \text{C.C.}$	$\mathbf{M}_e\delta\mathbf{R}_2(\psi_2) = \delta\mathbf{R}_2(\psi_2 + \sigma_2)$
(b) Unstable Confluent Res.	1 - Exp. Growth	$\delta\mathbf{R}_1 = A_1 e^{i\psi_1}\mathbf{E}_1 + \text{C.C.}$	$\mathbf{M}_e\delta\mathbf{R}_1(\psi_1) = \gamma_1\delta\mathbf{R}_1(\psi_1 + \sigma_1)$
	2 - Exp. Damping	$\delta\mathbf{R}_2 = A_2 e^{i\psi_2}\mathbf{E}_2 + \text{C.C.}$	$\mathbf{M}_e\delta\mathbf{R}_2(\psi_2) = (1/\gamma_1)\delta\mathbf{R}_2(\psi_2 + \sigma_1)$
(c) Unstable Lattice Res.	1 - Stable Osc.	$\delta\mathbf{R}_1 = A_1 e^{i\psi_1}\mathbf{E}_1 + \text{C.C.}$	$\mathbf{M}_e\delta\mathbf{R}_1(\psi_1) = \delta\mathbf{R}_1(\psi_1 + \sigma_1)$
	2 - Exp. Growth	$\delta\mathbf{R}_2 = A_2\mathbf{E}_2$	$\mathbf{M}_e\delta\mathbf{R}_2 = -\gamma_2\delta\mathbf{R}_2$
	3 - Exp. Damping	$\delta\mathbf{R}_3 = A_3\mathbf{E}_4$	$\mathbf{M}_e\delta\mathbf{R}_3 = -(1/\gamma_2)\delta\mathbf{R}_3$
(d) Unstable Double Lattice Resonance	1 - Exp. Growth	$\delta\mathbf{R}_1 = A_1\mathbf{E}_1$	$\mathbf{M}_e\delta\mathbf{R}_1 = -\gamma_1\delta\mathbf{R}_1$
	2 - Exp. Growth	$\delta\mathbf{R}_2 = A_2\mathbf{E}_2$	$\mathbf{M}_e\delta\mathbf{R}_2 = -\gamma_2\delta\mathbf{R}_2$
	3 - Exp. Damping	$\delta\mathbf{R}_3 = A_3\mathbf{E}_3$	$\mathbf{M}_e\delta\mathbf{R}_3 = -(1/\gamma_1)\delta\mathbf{R}_3$
	4 - Exp. Damping	$\delta\mathbf{R}_4 = A_4\mathbf{E}_4$	$\mathbf{M}_e\delta\mathbf{R}_4 = -(1/\gamma_2)\delta\mathbf{R}_4$

$$\delta\mathbf{R}_\ell \equiv \delta\mathbf{R}_\ell(s_i) \qquad \mathbf{E}_\ell \equiv \mathbf{E}_\ell(s_i) \qquad \mathbf{M}_e \equiv \mathbf{M}_e(s_i + L_p|s_i)$$

$$\delta\mathbf{R}(s) = \begin{cases} A_1[\mathbf{E}_1(s)e^{i\psi_1(s)} + \mathbf{E}_1^*(s)e^{-i\psi_1(s)}] + A_2[\mathbf{E}_2(s)e^{i\psi_2(s)} + \mathbf{E}_2^*(s)e^{-i\psi_2(s)}], & \text{cases (a) and (b)} \\ A_1[\mathbf{E}_1(s)e^{i\psi_1(s)} + \mathbf{E}_1^*(s)e^{-i\psi_1(s)}] + A_2\mathbf{E}_2(s) + A_3\mathbf{E}_4(s), & \text{case (c)} \\ A_1\mathbf{E}_1(s) + A_2\mathbf{E}_2(s) + A_3\mathbf{E}_3(s) + A_4\mathbf{E}_4(s), & \text{case (d)} \end{cases}$$

Decoupled Modes

In a continuous or periodic solenoidal focusing channel

$$\kappa_x(s) = \kappa_y(s) = \kappa(s)$$

with a round matched-beam solution

$$\varepsilon_x = \varepsilon_y \equiv \varepsilon = \text{const}$$
$$r_{xm}(s) = r_{ym}(s) \equiv r_m(s)$$

envelope perturbations are simply decoupled with:

Breathing Mode: $$\delta r_+ \equiv \frac{\delta r_x + \delta r_y}{2}$$

Quadrupole Mode: $$\delta r_- \equiv \frac{\delta r_x - \delta r_y}{2}$$

The resulting decoupled envelope equations are:

Breathing Mode:

$$\delta r''_+ + \kappa\delta r_+ + \frac{Q}{r_m^2}\delta r_+ + \frac{3\varepsilon^2}{r_m^4}\delta r_+ = -r_m\left(\frac{\delta\kappa_x + \delta\kappa_y}{2}\right) + \frac{1}{r_m}\delta Q + \frac{2\varepsilon}{r_m^3}\left(\frac{\delta\varepsilon_x + \delta\varepsilon_y}{2}\right)$$

Quadrupole Mode:

$$\delta r''_- + \kappa\delta r_- + \frac{3\varepsilon^2}{r_m^4}\delta r_- = -r_m\left(\frac{\delta\kappa_x - \delta\kappa_y}{2}\right) + \frac{2\varepsilon}{r_m^3}\left(\frac{\delta\varepsilon_x - \delta\varepsilon_y}{2}\right)$$

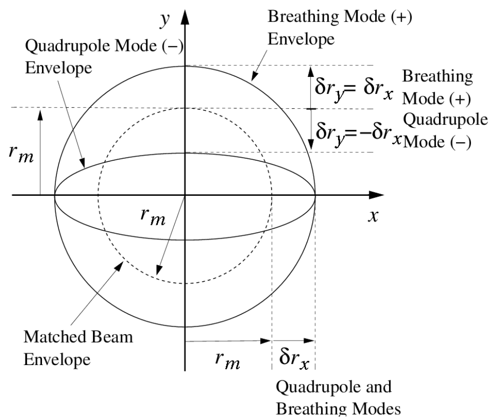
Graphical interpretation of mode symmetries:

Breathing Mode:

$$\delta r_+ = \frac{\delta r_x + \delta r_y}{2}$$

Quadrupole Mode:

$$\delta r_- = \frac{\delta r_x - \delta r_y}{2}$$



Decoupled Mode Properties:

Space charge terms ~ Q only directly expressed in equation for $\delta r_+(s)$
- Indirectly present in both equations from matched envelope $r_m(s)$

Homogeneous Solution:
- Restoring term for $\delta r_+(s)$ larger than for $\delta r_-(s)$
  - Breathing mode should oscillate faster than the quadrupole mode

Particular Solution:
- Misbalances in focusing and emittance driving terms can project onto either mode
  - nonzero perturbed $\kappa_x(s)$ + $\kappa_y(s)$ and $\varepsilon_x(s)$ + $\varepsilon_y(s)$ project onto breathing mode
  - nonzero perturbed $\kappa_x(s)$ - $\kappa_y(s)$ and $\varepsilon_x(s)$ - $\varepsilon_y(s)$ project onto quadrupole mode
- Perveance driving perturbations project *only* on breathing mode

Previous symmetry classes greatly reduce for decoupled modes:

Previous homogeneous 4x4 solution map:

$$\delta\mathbf{R}(s) = \mathbf{M}_e(s|s_i)\cdot\delta\mathbf{R}(s_i)$$
$$\delta\mathbf{R}(s) = (\delta r_x, \delta r_x', \delta r_y, \delta r_y')$$
$$\mathbf{M}_e(s|s_i) = \; 4\times 4 \text{ transfer map}$$

reduces to two independent 2x2 maps with greatly simplified symmetries:

$$\delta\mathbf{R} \equiv (\delta r_+, \delta r_+', \delta r_-, \delta r_-')$$

$$\mathbf{M}_e(s_i+L_p|s_i) = \begin{bmatrix} \mathbf{M}_+(s_i+L_p|s_i) & 0 \\ 0 & \mathbf{M}_-(s_i+L_p|s_i) \end{bmatrix}$$

with corresponding eigenvalue problems:

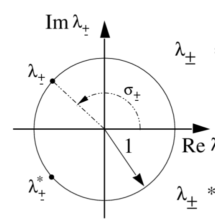
$$\mathbf{M}_\pm(s_i+L_p|s_i)\cdot\mathbf{E}_n(s_i) = \lambda_\pm \mathbf{E}_n(s_i)$$

Many familiar results from analysis of Hills equation (see: S.M. Lund lectures on Transverse Particle Dynamics) can be immediately applied to the decoupled case, for example:

$$\frac{1}{2}|\mathrm{Tr}\;\mathbf{M}_\pm(s_i+L_p|s_i)| \leq 1 \quad \Longleftrightarrow \quad \text{mode stability}$$

Eigenvalue symmetries and launching conditions simplify for decoupled modes

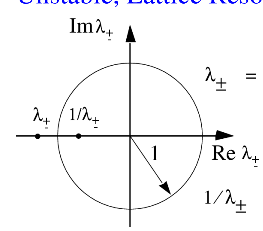
Eigenvalue Symmetry 1:
Stable



$\lambda_\pm = e^{i\sigma_\pm}$

$\lambda_\pm^* = 1/\lambda_\pm = e^{-i\sigma_\pm}$

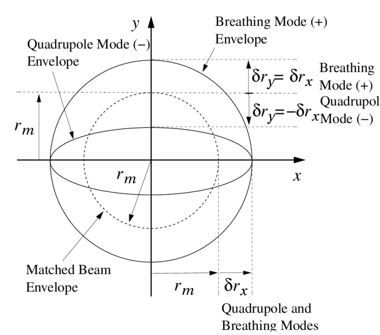
Eigenvalue Symmetry 2:
Unstable, Lattice Resonance



$\lambda_\pm = \gamma_\pm e^{-i\pi}$

$1/\lambda_\pm = (1/\gamma_\pm)e^{-i\pi}$

Launching Condition / Projections



General Mode Limits

Using phase-amplitude analysis can show for any linear focusing lattice:

1) Phase advance of any normal mode satisfies the zero space-charge limit:

$$\lim_{Q\to 0} \sigma_\ell = 2\sigma_0$$

2) Pure normal modes (not driven) evolve with a quadratic phase-space (Courant-Snyder) invariant in the normal coordinates of the mode

Simply expressed for decoupled modes with $\kappa_x = \kappa_y, \;\; \varepsilon_x = \varepsilon_y$

$$\left[\frac{\delta r_\pm(s)}{w_\pm(s)}\right]^2 + [w'_\pm(s)\delta r_\pm(s) - w_\pm(s)\delta r'_\pm(s)]^2 = \text{const}$$

where

$$w''_+ + \kappa w_+ + \frac{Q}{r_m^2} w_+ + \frac{3\varepsilon^2}{r_m^4} w_+ - \frac{1}{w_+^3} = 0$$

$$w''_- + \kappa w_- + \frac{3\varepsilon^2}{r_m^4} w_- - \frac{1}{w_-^3} = 0$$

$$w_\pm(s + L_p) = w_\pm(s)$$

Analogous results for coupled modes [See Edwards and Teng, IEEE Trans Nuc. Sci. **20**, 885 (1973)]

- More complex expression due to coupling

S7: Envelope Modes in Continuous Focusing

Lund and Bukh, PRSTAB **7**, 024801 (2004)

Focusing: $\kappa_x(s) = \kappa_y(s) = k_{\beta 0}^2 = \left(\frac{\sigma_0}{L_p}\right)^2 = \text{const}$

Matched beam: $\varepsilon_x = \varepsilon_y = \varepsilon = \text{const}$

symmetric beam: $r_{xm}(s) = r_{ym}(s) = r_m = \text{const}$

matched envelope: $k_{\beta 0}^2 r_m - \frac{Q}{r_m} - \frac{\varepsilon^2}{r_m^3} = 0$

depressed phase advance: $\sigma = \sqrt{\sigma_0^2 - \frac{Q}{(r_m/L_p)^2}} = \frac{\varepsilon L_p}{r_m^2}$

one parameter needed for scaled solution: $\frac{k_{\beta 0}^2 \varepsilon^2}{Q^2} = \frac{\sigma_0^2 \varepsilon^2}{Q^2 L_p^2} = \frac{(\sigma/\sigma_0)^2}{[1 - (\sigma/\sigma_0)^2]^2}$

Decoupled Modes:

$$\delta r_\pm(s) = \frac{\delta r_x(s) \pm \delta r_y(s)}{2}$$

Envelope equations of motion become:

$$L_p^2 \frac{d^2}{ds^2}\left(\frac{\delta r_+}{r_m}\right) + \sigma_+^2 \left(\frac{\delta r_+}{r_m}\right) = -\frac{\sigma_0^2}{2}\left(\frac{\delta\kappa_x}{k_{\beta 0}^2} + \frac{\delta\kappa_y}{k_{\beta 0}^2}\right) + (\sigma_0^2 - \sigma^2)\frac{\delta Q}{Q} + \sigma^2\left(\frac{\delta\varepsilon_x}{\varepsilon} + \frac{\delta\varepsilon_y}{\varepsilon}\right)$$

$$L_p^2 \frac{d^2}{ds^2}\left(\frac{\delta r_-}{r_m}\right) + \sigma_-^2 \left(\frac{\delta r_-}{r_m}\right) = -\frac{\sigma_0^2}{2}\left(\frac{\delta\kappa_x}{k_{\beta 0}^2} - \frac{\delta\kappa_y}{k_{\beta 0}^2}\right) + \sigma^2\left(\frac{\delta\varepsilon_x}{\varepsilon} - \frac{\delta\varepsilon_y}{\varepsilon}\right)$$

$\sigma_+ \equiv \sqrt{2\sigma_0^2 + 2\sigma^2}$ "breathing" mode phase advance

$\sigma_- \equiv \sqrt{\sigma_0^2 + 3\sigma^2}$ "quadrupole" mode phase advance

Homogeneous equations for normal modes:

$$\frac{d^2}{ds^2}\delta r_\pm + \left(\frac{\sigma_\pm}{L_p}\right)^2 \delta r_\pm = 0$$

See also lectures by J.J. Barnard, Envelope Modes and Halo

- Simple harmonic oscillator equation

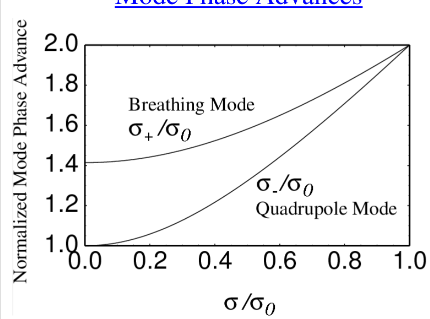
Homogeneous Solution (normal modes):

$$\delta r_\pm(s) = \delta r_\pm(s_i)\cos\left(\sigma_\pm \frac{s - s_i}{L_p}\right) + \frac{\delta r'_\pm(s_i)}{\sigma_\pm/L_p}\sin\left(\sigma_\pm \frac{s - s_i}{L_p}\right)$$

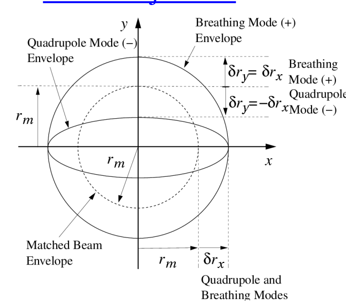
$\delta r_\pm(s_i), \;\; \delta r'_\pm(s_i)$ mode initial conditions

Properties of continuous focusing homogeneous solution: Normal Modes

Mode Phase Advances



Mode Projections



$\sigma_+ \equiv \sqrt{2\sigma_0^2 + 2\sigma^2}$ Breathing Mode: $\delta r_+ \equiv \frac{\delta r_x + \delta r_y}{2}$

$\sigma_- \equiv \sqrt{\sigma_0^2 + 3\sigma^2}$ Quadrupole Mode: $\delta r_- \equiv \frac{\delta r_x - \delta r_y}{2}$

Particular Solution (driving perturbations):

Green's function form of solution derived using projections onto normal modes

- See proof that this is a valid solution is given in Appendix A

$$\frac{\delta r_\pm(s)}{r_m} = \frac{1}{L_p^2}\int_{s_i}^{s} d\tilde{s}\; G_\pm(s,\tilde{s})\delta p_\pm(\tilde{s})$$

$$\delta p_+(s) = -\frac{\sigma_0^2}{2}\left[\frac{\delta\kappa_x(s)}{k_{\beta0}^2} + \frac{\delta\kappa_y(s)}{k_{\beta0}^2}\right] + (\sigma_0^2 - \sigma^2)\frac{\delta Q(s)}{Q} + \sigma^2\left[\frac{\delta\varepsilon_x(s)}{\varepsilon} + \frac{\delta\varepsilon_y(s)}{\varepsilon}\right]$$

$$\delta p_-(s) = -\frac{\sigma_0^2}{2}\left[\frac{\delta\kappa_x(s)}{k_{\beta0}^2} - \frac{\delta\kappa_y(s)}{k_{\beta0}^2}\right] + \sigma^2\left[\frac{\delta\varepsilon_x(s)}{\varepsilon} - \frac{\delta\varepsilon_y(s)}{\varepsilon}\right]$$

$$G_\pm(s,\tilde{s}) = \frac{1}{\sigma_\pm/L_p}\sin\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right)$$

Green's function solution is *fully general*. Insight gained from simplified solutions for specific classes of driving perturbations:

- Adiabatic
- Sudden

covered in these lectures

- Ramped
- Harmonic

covered in PRSTAB Review article

Continuous Focusing – adiabatic particular solution

For driving perturbations $\delta p_+(s)$ and $\delta p_-(s)$ slow on quadrupole mode (slower mode) wavelength $\sim 2\pi L_p/\sigma_-$ the Green function solution reduces to:

$$\frac{\delta r_+(s)}{r_m} = \frac{\delta p_+(s)}{\sigma_+^2}$$

$$= -\left[\frac{1}{2}\frac{1}{1+(\sigma/\sigma_0)^2}\right]\frac{1}{2}\left(\frac{\delta\kappa_x(s)}{k_{\beta0}^2} + \frac{\delta\kappa_y(s)}{k_{\beta0}^2}\right) + \left[\frac{1}{2}\frac{1-(\sigma/\sigma_0)^2}{1+(\sigma/\sigma_0)^2}\right]\frac{\delta Q(s)}{Q}$$

$$+ \left[\frac{(\sigma/\sigma_0)^2}{1+(\sigma/\sigma_0)^2}\right]\frac{1}{2}\left(\frac{\delta\varepsilon_x(s)}{\varepsilon} + \frac{\delta\varepsilon_y(s)}{\varepsilon}\right),$$

Focusing; Perveance; Emittance

Coefficients of adiabatic terms in square brackets"[ ]"

$$\frac{\delta r_-(s)}{r_m} = \frac{\delta p_-(s)}{\sigma_-^2}$$

$$= -\left[\frac{1}{1+3(\sigma/\sigma_0)^2}\right]\frac{1}{2}\left(\frac{\delta\kappa_x(s)}{k_{\beta0}^2} - \frac{\delta\kappa_y(s)}{k_{\beta0}^2}\right)$$

$$+ \left[\frac{2(\sigma/\sigma_0)^2}{1+3(\sigma/\sigma_0)^2}\right]\frac{1}{2}\left(\frac{\delta\varepsilon_x(s)}{\varepsilon} - \frac{\delta\varepsilon_y(s)}{\varepsilon}\right).$$

Focusing; Emittance

Derivation of Adiabatic Solution:

- Several ways to derive, show more "mechanical" procedure here ....

Use:

$$\frac{\delta r_\pm(s)}{r_m} = \frac{1}{L_p^2}\int_{s_i}^{s} d\tilde{s}\; G_\pm(s,\tilde{s})\delta p_\pm(\tilde{s})$$

$$G_\pm(s,\tilde{s}) = \frac{1}{\sigma_\pm/L_p}\sin\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right) = \frac{1}{(\sigma_\pm/L_p)^2}\frac{d}{d\tilde{s}}\cos\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right)$$

Gives:

$$\frac{\delta r_\pm(s)}{r_m} = \int_{s_i}^{s} d\tilde{s}\left[\frac{d}{d\tilde{s}}\cos\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right)\right]\frac{\delta p_\pm(\tilde{s})}{\sigma_\pm}$$

$$= \int_{s_i}^{s} d\tilde{s}\,\frac{d}{d\tilde{s}}\left[\cos\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right)\frac{\delta p_\pm(\tilde{s})}{\sigma_\pm}\right] - \int_{s_i}^{s} d\tilde{s}\,\cos\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right)\frac{d}{d\tilde{s}}\frac{\delta p_\pm(\tilde{s})}{\sigma_\pm}$$

$$= \cos\left(\sigma_\pm\frac{s-\tilde{s}}{L_p}\right)\frac{\delta p_\pm(\tilde{s})}{\sigma_\pm}\Bigg|_{\tilde{s}=s_i}^{\tilde{s}=s} = \frac{\delta p_\pm(s)}{\sigma_\pm} - \cos\left(\sigma_\pm\frac{s-s_i}{L_p}\right)\frac{\delta p_\pm(s_i)}{\sigma_\pm}$$

$$= \frac{\delta p_\pm(s)}{\sigma_\pm}$$

Adiabatic → 0

No Initial Perturbation → 0

Comments on Adiabatic Solution:

- Adiabatic response is essentially a slow adaptation in the matched envelope to perturbations (solution does not oscillate due to slow changes)
- Slow envelope frequency σ_- sets the scale for slow variations required

Replacements in adiabatically adapted match:

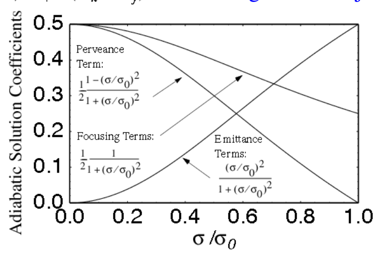
$$r_x = r_m \rightarrow r_m + \delta r_+ + \delta r_-$$
$$r_y = r_m \rightarrow r_m + \delta r_- - \delta r_+$$

Parameter replacements in rematched beam (no longer axisymmetric):

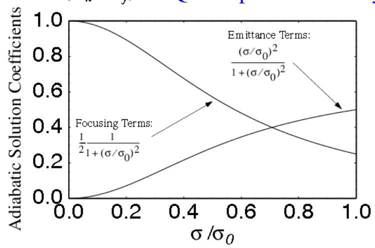
$$\kappa_x = k_{\beta0}^2 \rightarrow k_{\beta0}^2 + \delta\kappa_x(s)$$
$$\kappa_y = k_{\beta0}^2 \rightarrow k_{\beta0}^2 + \delta\kappa_y(s)$$
$$Q \rightarrow Q + \delta Q(s)$$
$$\varepsilon_x = \varepsilon \rightarrow \varepsilon + \delta\varepsilon_x(s)$$
$$\varepsilon_y = \varepsilon \rightarrow \varepsilon + \delta\varepsilon_y(s)$$

Continuous Focusing – adiabatic solution coefficients

a) $\delta r_+ = (\delta r_x + \delta r_y)/2$ Breathing Mode Projection



b) $\delta r_- = (\delta r_x - \delta r_y)/2$ Quadrupole Mode Projection



Relative strength of:

- Space-Charge (Perveance)
- Applied Focusing
- Emittance

terms vary with space-charge depression (σ/σ_0) for both breathing and quadrupole mode projections

Plots allow one to read off the relative importance of various contributions to beam mismatch as a function of space-charge strength

Continuous Focusing – sudden particular solution

For sudden, step function driving perturbations of form:

$$\delta p_\pm(s) = \widehat{\delta p_\pm}\Theta(s - s_p)$$

$s = s_p =$ axial coordinate perturbation applied

Hat quantities are constant amplitudes

with amplitudes:

$$\widehat{\delta p_+} = -\frac{\sigma_0^2}{2}\left[\frac{\widehat{\delta\kappa_x}}{k_{\beta 0}^2} + \frac{\widehat{\delta\kappa_y}}{k_{\beta 0}^2}\right] + (\sigma_0^2 - \sigma^2)\frac{\widehat{\delta Q}}{Q} + \sigma^2\left[\frac{\widehat{\delta\varepsilon_x}}{\varepsilon} + \frac{\widehat{\delta\varepsilon_y}}{\varepsilon}\right] = \text{const}$$

$$\widehat{\delta p_-} = -\frac{\sigma_0^2}{2}\left[\frac{\widehat{\delta\kappa_x}}{k_{\beta 0}^2} - \frac{\widehat{\delta\kappa_y}}{k_{\beta 0}^2}\right] + \sigma^2\left[\frac{\widehat{\delta\varepsilon_x}}{\varepsilon} - \frac{\widehat{\delta\varepsilon_y}}{\varepsilon}\right] = \text{const}$$

The solution is given by the substitution in the expression for the adiabatic solution:

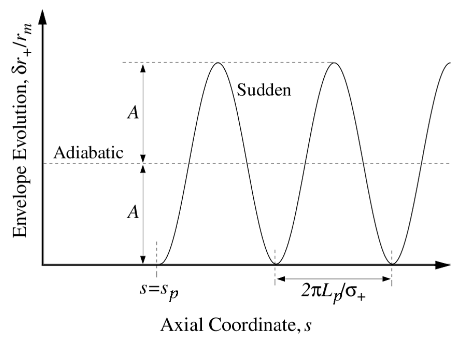
- Manipulate Green's function solution to show (similar to Adiabatic case steps)

$$\delta p_\pm(s) \to \widehat{\delta p_\pm}\left[1 - \cos\left(\sigma_\pm \frac{s - s_p}{L_p}\right)\right]\Theta(s - s_p)$$

Sudden perturbation solution, substitute in pervious adiabatic expressions:

$$\delta p_\pm(s) \to \widehat{\delta p_\pm}\left[1 - \cos\left(\sigma_\pm \frac{s - s_p}{L_p}\right)\right]\Theta(s - s_p)$$

Illustration of solution properties for a sudden $\delta p_+(s)$ perturbation term



For the same amplitude of total driving perturbations, sudden perturbations result in 2x the envelope excursion that adiabatic perturbations produce

Continuous Focusing – Driven perturbations on a continuously focused matched equilibrium (summary)

Adiabatic Perturbations:

- Essentially a rematch of equilibrium beam if the change is slow relative to quadrupole envelope mode oscillations

Sudden Perturbations:

- Projects onto breathing and quadrupole envelope modes with 2x adiabatic amplitude oscillating from zero to max amplitude

Ramped Perturbations: (see PRSTAB article; based on Green's function)

- Can be viewed as a superposition between the adiabatic and sudden form perturbations

Harmonic Perturbations: (see PRSTAB article; based on Green's function)

- Can build very general cases of driven perturbations by linear superposition
- Results may be less "intuitive" (expressed in complex form)

Cases covered in class illustrate a range of common behavior and help build intuition on what can drive envelope oscillations and the relative importance of various terms as a function of space-charge strength

Appendix A: Particular Solution for Driven Envelope Modes

Lund and Bukh, PRSTAB **7**, 024801 (2004)

Following Wiedemann (Particle Accelerator Physics, 1993, pp 106) first, consider more general *Driven Hill's Equation*

$$x'' + \kappa(s)x = p(s)$$

The corresponding homogeneous equation:

$$x'' + \kappa(s)x = 0$$

has principal solutions

$$x(s) = C_1\mathcal{C}(s) + C_2\mathcal{S}(s) \qquad C_1, C_2 = \text{constants}$$

where

Cosine-Like Solution	Sine-Like Solution
$\mathcal{C}'' + \kappa(s)\mathcal{C} = 0$	$\mathcal{S}'' + \kappa(s)\mathcal{S} = 0$
$\mathcal{C}(s = s_i) = 1$	$\mathcal{S}(s = s_i) = 0$
$\mathcal{C}'(s = s_i) = 0$	$\mathcal{S}'(s = s_i) = 1$

Recall that the homogeneous solutions have the Wronskian symmetry:

- See S.M. Lund lectures on Transverse Dynamics, S5C

$$W(s) = \mathcal{C}(s)\mathcal{S}'(s) - \mathcal{C}'(s)\mathcal{S}(s) = 1$$

A particular solution to the *Driven Hill's Equation* can be constructed using a Greens' function method:

$$x(s) = \int_{s_i}^{s} d\tilde{s}\; G(s,\tilde{s})p(\tilde{s})$$

$$G(s,\tilde{s}) = \mathcal{S}(s)\mathcal{C}(\tilde{s}) - \mathcal{C}(s)\mathcal{S}(\tilde{s})$$

Demonstrate this works by first taking derivatives:

$$x = \mathcal{S}(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{C}(\tilde{s})p(\tilde{s}) \;-\; \mathcal{C}(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{S}(\tilde{s})p(\tilde{s})$$

$$x' = \mathcal{S}'(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{C}(\tilde{s})p(\tilde{s}) \;-\; \mathcal{C}'(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{S}(\tilde{s})p(\tilde{s}) + p(s)\underbrace{[\mathcal{S}(s)\mathcal{C}(s) - \mathcal{S}(s)\mathcal{C}(s)]}_{0}$$

$$= \mathcal{S}'(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{C}(\tilde{s})p(\tilde{s}) \;-\; \mathcal{C}'(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{S}(\tilde{s})p(\tilde{s})$$

$$x'' = \mathcal{S}''(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{C}(\tilde{s})p(\tilde{s}) \;-\; \mathcal{C}''(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{S}(\tilde{s})p(\tilde{s}) + p(s)\underbrace{[\mathcal{S}'(s)\mathcal{C}(s) - \mathcal{C}'(s)\mathcal{S}(s)]}_{1}$$

Wronskian Symmetry

$$= p(s) \;+\; \mathcal{S}''(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{C}(\tilde{s})p(\tilde{s}) \;-\; \mathcal{C}''(s)\int_{s_i}^{s} d\tilde{s}\;\mathcal{S}(\tilde{s})p(\tilde{s})$$

Apply these results in the *Driven Hill's Equation*:

Definition of Principal Orbit Functions

$$x'' + \kappa(s)x = p(s) \;+\; \underbrace{[\mathcal{S}'' + \kappa\mathcal{S}]}_{0}\int_{s_i}^{s} d\tilde{s}\;\mathcal{C}(\tilde{s})p(\tilde{s}) \;-\; \underbrace{[\mathcal{C}'' + \kappa\mathcal{C}]}_{0}\int_{s_i}^{s} d\tilde{s}\;\mathcal{S}(\tilde{s})p(\tilde{s})$$

$$= p(s)$$

Thereby proving we have a valid particular solution. The general solution to the *Driven Hill's Equation* is then:

- Choose constants $C_1,\ C_2$ consistent with particle initial conditions at $s = s_i$

$$x(s) = x(s_i)\mathcal{C}(s) \;+\; x'(s_i)\mathcal{S}(s) \;+\; \int_{s_i}^{s} d\tilde{s}\; G(s,\tilde{s})p(\tilde{s})$$

$$G(s,\tilde{s}) = \mathcal{S}(s)\mathcal{C}(\tilde{s}) - \mathcal{C}(s)\mathcal{S}(\tilde{s})$$

Apply these results to the driven perturbed envelope equation:

$$\frac{d^2}{ds^2}\delta r_\pm + \frac{\sigma_\pm^2}{L_p^2}\delta r_\pm = \frac{r_m}{L_p^2}\delta p_\pm$$

The homogeneous equations can be solved exactly for continuous focusing:

$$\mathcal{C}(s) = \cos\left(\sigma_\pm\frac{s - s_i}{L_p}\right)$$

$$\mathcal{S}(s) = \frac{L_p}{\sigma_\pm}\sin\left(\sigma_\pm\frac{s - s_i}{L_p}\right)$$

and the Green's function can be simplified as:

$$G(s,\tilde{s}) = \mathcal{S}(s)\mathcal{C}(\tilde{s}) - \mathcal{C}(s)\mathcal{S}(\tilde{s})$$

$$= \frac{L_p}{\sigma_\pm}\left\{\sin\left(\sigma_\pm\frac{s - s_i}{L_p}\right)\cos\left(\sigma_\pm\frac{\tilde{s} - s_i}{L_p}\right) - \cos\left(\sigma_\pm\frac{s - s_i}{L_p}\right)\sin\left(\sigma_\pm\frac{\tilde{s} - s_i}{L_p}\right)\right\}$$

$$= \frac{L_p}{\sigma_\pm}\sin\left(\sigma_\pm\frac{s - \tilde{s}}{L_p}\right)$$

Using these results the particular solution for the driven perturbed envelope equation can be expressed as:

- Here we rescale the Green's function to put in the form given in S8

$$\frac{\delta r_\pm(s)}{r_m} = \frac{1}{L_p^2}\int_{s_i}^{s} d\tilde{s}\; G_\pm(s,\tilde{s})\delta p_\pm(\tilde{s})$$

$$G_\pm(s,\tilde{s}) = \frac{1}{\sigma_\pm/L_p}\sin\left(\sigma_\pm\frac{s - \tilde{s}}{L_p}\right)$$

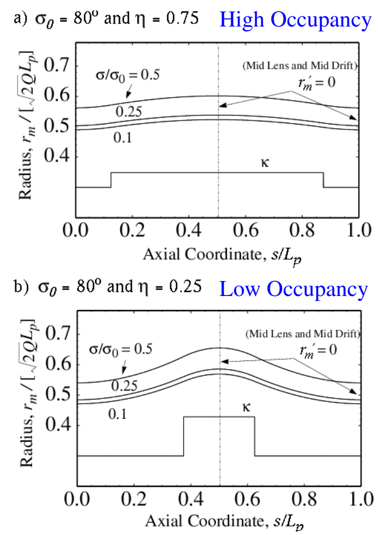
S8: Envelope Modes in Periodic Focusing Channels

Lund and Bukh, PRSTAB **7**, 024801 (2004)

Overview

- Much more complicated than continuous focusing results
  - Lattice can couple to oscillations and destabilize the system
  - Broad parametric instability bands can result
- Instability bands calculated will exclude wide ranges of parameter space from machine operation
  - Exclusion region depends on focusing type
  - Will find that alternating gradient quadrupole focusing tends to have more instability than high occupancy solenoidal focusing due to larger envelope flutter driving stronger, broader instability
- Results in this section are calculated numerically and summarized parametrically to illustrate the full range of normal mode characteristics
  - Driven modes not considered but should be mostly analogous to CF case
  - Results presented in terms of phase advances and normalized space-charge strength to allow broad applicability
  - Coupled 4x4 eigenvalue problem and mode symmetries identified in S6 are solved numerically and analytical limits are verified
  - Carried out for piecewise constant lattices for simplicity (fringe changes little)
- More information on results presented can be found in the PRSTAB review

Solenoidal Focusing – Matched Envelope Solution



Focusing:

$$\kappa_x(s) = \kappa_y(s) = \kappa(s)$$

$$\kappa(s + L_p) = \kappa(s)$$

Matched Beam:

$$\varepsilon_x = \varepsilon_y = \varepsilon = \text{const}$$

$$r_{xm}(s) = r_{ym}(s) = r_m(s)$$

$$r_m(s + L_p) = r_m(s)$$

Comments:

- Envelope flutter a strong function of occupancy η
- Space-charge expands envelope but does not strongly modify periodic flutter

Envelope Flutter Scaling of Matched Envelope Solution

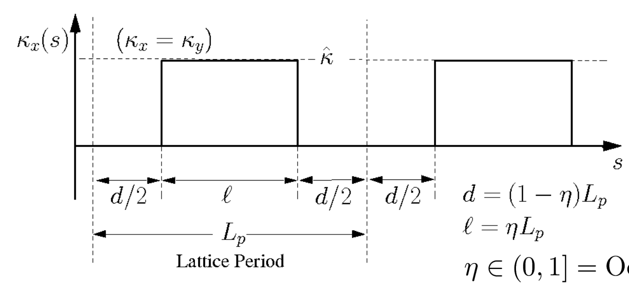
Add material explaining scaling better in future editions

Using a transfer matrix approach on undepressed single-particle orbits set the strength of the focusing function for specified undepressed particle phase advance by solving:

- See: S.M. Lund, lectures on Transverse Particle Dynamics
- Particle phase-advance is measured in the rotating Larmor frame

Solenoidal Focusing - piecewise constant focusing lattice

$$\cos\sigma_0 = \cos(2\Theta) - \frac{1-\eta}{\eta}\Theta\sin(2\Theta) \qquad \Theta \equiv \frac{\sqrt{\hat{\kappa}}L_p}{2}$$

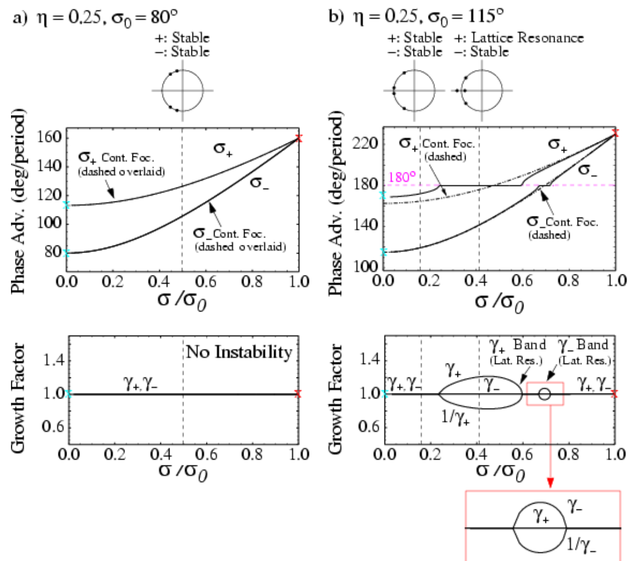


$d = (1-\eta)L_p$

$\ell = \eta L_p$

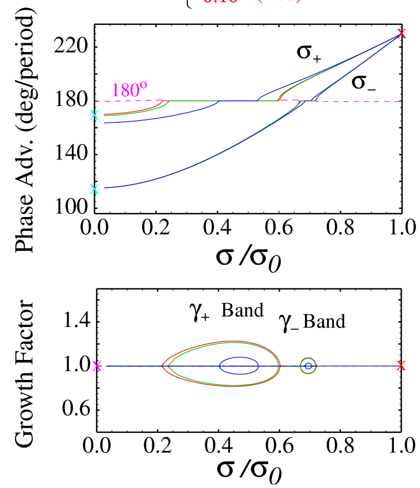
$\eta \in (0, 1] = \text{Occupancy}$

Solenoidal Focusing – parametric plots of breathing and quadrupole envelope mode phase advances two values of undepressed phase advance



Solenoidal Focusing – mode instability bands become wider and stronger for smaller occupancy

$$\eta = \begin{cases} 0.75 & \text{(Blue)} \\ 0.25 & \text{(Green)} \\ 0.10 & \text{(Red)} \end{cases} \qquad \sigma_0 = 115^\circ$$

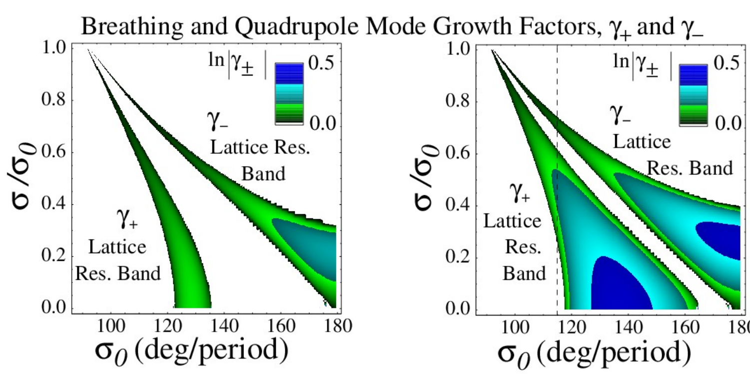


Comments:

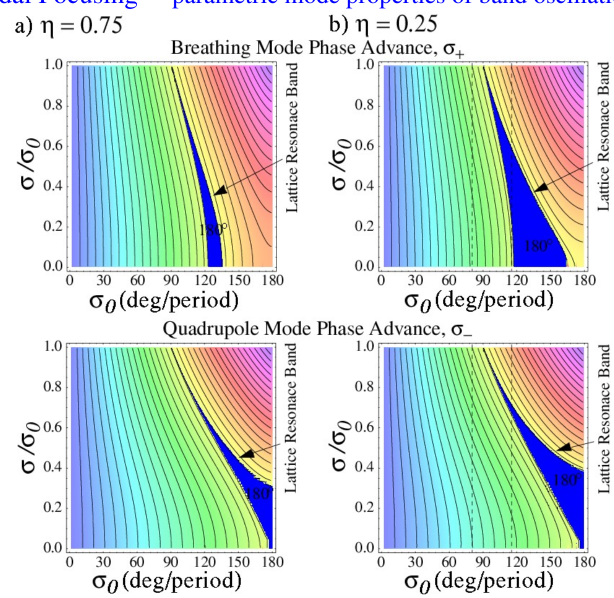
- Mode phase advance in instability band 180 degrees per lattice period
- Significant deviations from continuous model even outside the band of instability when space-charge is strong
- Instability band becomes stronger/broader for low occupancy and weaker/narrower for high occupancy
  - Disappears at full occupancy (continuous limit)

Solenoidal Focusing – broad ranges of parametric instability are found for the breathing and quadrupole bands that must be avoided in machine operation

$\eta = 0.75$ $\eta = 0.25$



Solenoidal Focusing – parametric mode properties of band oscillations



Parametric scaling of the boundary of the region of instability

Solenoid instability bands identified as a Lattice Resonance Instability corresponding to a 1/2-integer parametric resonance between the mode oscillation frequency and the lattice

Estimate normal mode frequencies for weak focusing from continuous focusing theory:

$$\sigma_+ \simeq \sqrt{2\sigma_0^2 + 2\sigma^2}$$

$$\sigma_- \simeq \sqrt{\sigma_0^2 + 3\sigma^2}$$

This gives (measure phase advance in degrees):

<u>Breathing Band</u>:

$$\sigma_+ = 180^\circ$$

$$\Longrightarrow \quad \sqrt{2\sigma_0^2 + 2\sigma^2} = 180^\circ$$

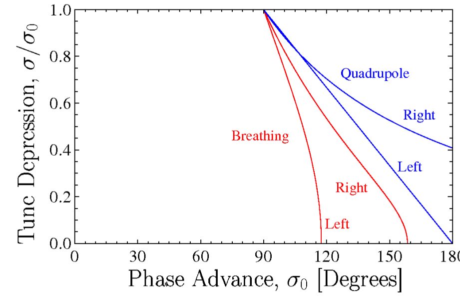
<u>Quadrupole Band</u>:

$$\sigma_- = 180^\circ$$

$$\Longrightarrow \quad \sqrt{\sigma_0^2 + 3\sigma^2} = 180^\circ$$

- Predictions poor due to inaccurate mode frequency estimates
  - Predictions nearer to left edge of band rather than center (expect resonance strongest at center)
- Simple resonance condition cannot predict width of band
  - Important to characterize width to avoid instability in machine designs
  - Width of band should vary strongly with solenoid occupancy η

To provide a practical guide on the location/width of the breathing and quadrupole envelope bands, many parametric runs were made and the instability band boundaries were quantified through curve fitting:



<u>Breathing Band Boundaries</u>:

$$\sigma^2 + f\sigma_0^2 = (90^\circ)^2(1+f)$$

$$f = f(\sigma_0, \eta) = \begin{cases} 1.113 - 0.413\eta + 0.00348\sigma_0, & \text{left-edge} \\ 1.046 + 0.318\eta - 0.00410\sigma_0, & \text{right-edge} \end{cases}$$

<u>Quadrupole Band Boundaries</u>:

Left: $\sigma/\sigma_0 + g\dfrac{\sigma_0}{90^\circ} = 1 + g$

Right: $\sigma + g\sigma_0 = 90^\circ(1+g)$

$$g = g(\eta) = \begin{cases} 1, & \text{left-edge} \\ 0.227 - 0.173\eta, & \text{right-edge} \end{cases}$$

- Breathing band: maximum errors ~5 /~2 degrees on left/right boundaries
- Quadrupole band: maximum errors ~8/~3 degrees on left/right boundaries

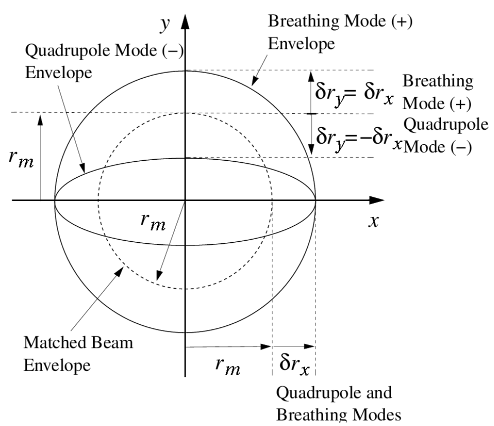
Pure eigenmode launching conditions are simple for the ideal solenoid case and correspond to the breathing (+) and quadrupole (-) mode symmetries covered for decoupoled modes in S6

Breathing Mode:

$$\delta r_+ = \frac{\delta r_x + \delta r_y}{2}$$

Quadrupole Mode:

$$\delta r_- = \frac{\delta r_x - \delta r_y}{2}$$

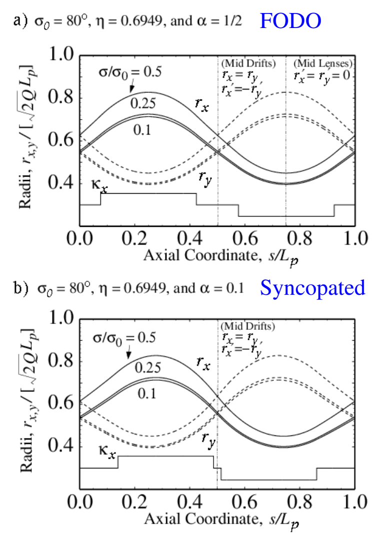


Caution:

Recall we are describing problem implicitly in the rotating (Larmor) frame and to express launch conditions in the lab frame quadrupole mode conditions must be projected back with the correct overall rotation through magnet fringe fields

Quadrupole Doublet Focusing – Matched Envelope Solution

FODO and Syncopated Lattices



Focusing:

$$\kappa_x(s) = -\kappa_y(s) = \kappa(s)$$

$$\kappa(s + L_p) = \kappa(s)$$

Matched Beam:

$$\varepsilon_x = \varepsilon_y = \varepsilon = \text{const}$$

$$r_{xm}(s + L_p) = r_{xm}(s)$$

$$r_{ym}(s + L_p) = r_{ym}(s)$$

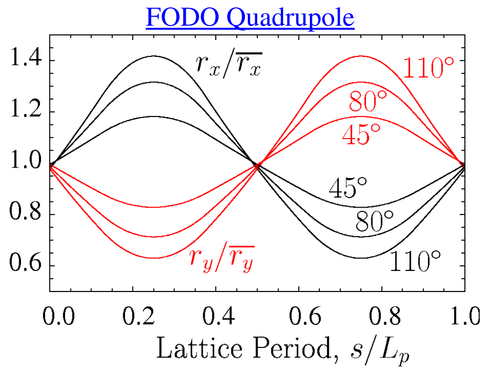
Comments:

- Envelope flutter a *weak* function of occupancy η
- Syncopation factors $\alpha \neq 1/2$ reduce envelope symmetry and can drive more instabilities
- Space-charge expands envelope

Envelope Flutter Scaling of Matched Envelope Solution

For FODO quadrupole transport, plot relative matched beam envelope excursions for a fixed form focusing lattice and fixed beam perveance as the strength of applied focusing strength increases as measured by σ_0

FODO Quadrupole



$$\overline{r_x} = \int_0^{L_p} \frac{ds}{L_p} r_x(s)$$

$\eta = 0.5 \quad L_p = 0.5$ m

$Q = 5 \times 10^{-4}$

$\varepsilon_x = \varepsilon_y = 50$ mm-mrad

σ_0	σ/σ_0
45°	0.20
80°	0.26
110°	0.32

- Larger matched envelope "flutter" corresponds to larger σ_0
  - More flutter results in higher prospects for instability due to transfer of energy from applied focusing
- Little dependence of flutter on quadrupole occupancy

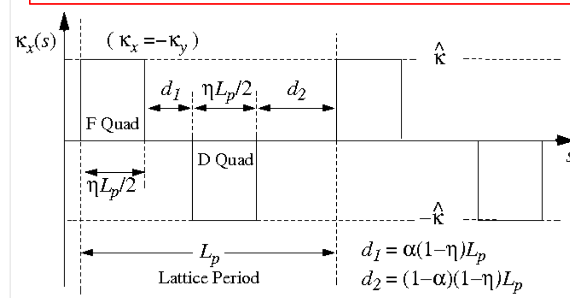
Using a transfer matrix approach on undepressed single-particle orbits set the strength of the focusing function for specified undepressed particle phase advance by solving:

- See: S.M. Lund, lectures on Transverse Particle Dynamics

Quadrupole Doublet Focusing - piecewise constant focusing lattice

$$\cos\sigma_0 = \cos\Theta\cosh\Theta + \frac{1-\eta}{\eta}\theta(\cos\Theta\sinh\Theta - \sin\Theta\cosh\Theta) - 2\alpha(1-\alpha)\frac{(1-\eta)^2}{\eta^2}\Theta^2\sin\Theta\sinh\Theta$$

$$\Theta \equiv \frac{\sqrt{|\hat{\kappa}|}L_p}{2}$$

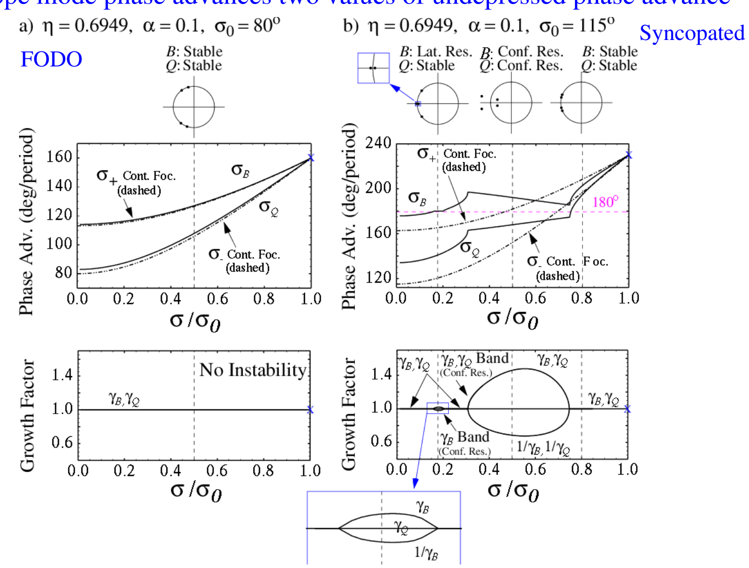


$\eta \in (0, 1]$ Occupancy

$\alpha \in [0, 1/2]$ Syncopation Factor

$\alpha = 1/2 \rightarrow$ FODO

Quadrupole Focusing – parametric plots of breathing and quadrupole envelope mode phase advances two values of undepressed phase advance



Important point:

For quadrupole focusing the normal mode coordinates are *NOT*

$$\delta r_\pm = \frac{\delta r_x \pm \delta r_y}{2}$$

$\delta r_+ \Leftrightarrow$ Breathing Mode

$\delta r_- \Leftrightarrow$ Quadrupole Mode

- Only works for axisymmetric focusing $(\kappa_x = \kappa_y = \kappa)$ with an axisymmetric matched beam $(\varepsilon_x = \varepsilon_y = \varepsilon)$

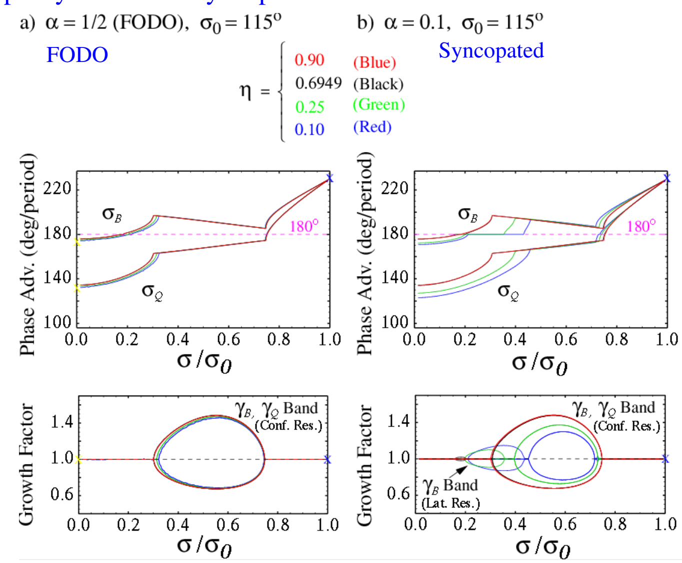
However, for low σ_0 we will find that the two stable modes correspond closely in frequency with continuous focusing model breathing and quadrupole modes even though they have different symmetry properties in terms of normal mode coordinates. Due to this, we denote:

Subscript B <==> Breathing Mode

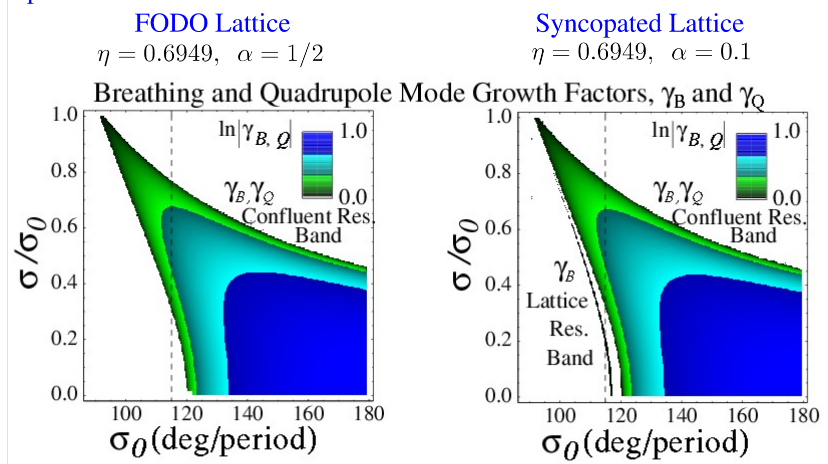
Subscript Q <==> Quadrupole Mode

- Label branches breathing and quadrupole in terms of low σ_0 branch frequencies corresponding to breathing and quadrupole frequencies from continuous theory
- Continue label to larger values of σ_0 where frequency correspondence with continuous modes breaks down

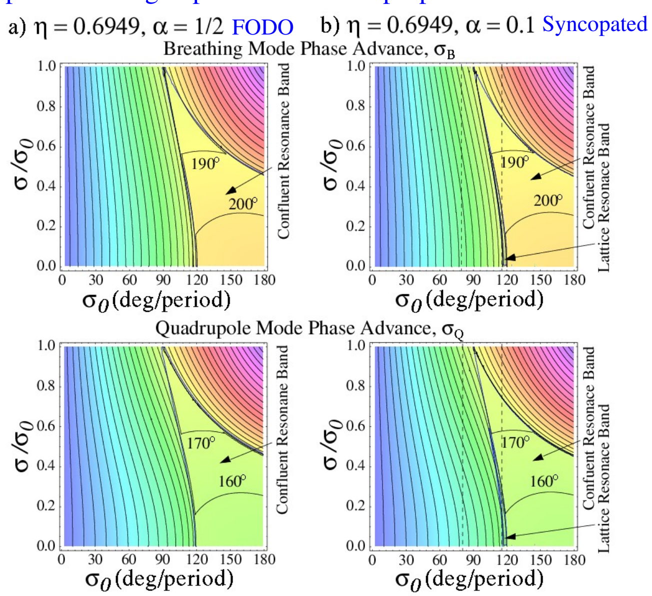
Quadrupole Focusing – mode instability bands vary little/strongly with occupancy for FODO/syncopated lattices



Quadrupole Focusing – broad ranges of parametric instability are found for the breathing and quadrupole bands that must be avoided in machine operation



Quadrupole Focusing – parametric mode properties of band oscillations



Parametric scaling of the boundary of the region of instability

Quadrupole instability bands identified:

- Confluent Band: 1/2-integer parametric resonance between *both* breathing and quadrupole modes and the lattice
- Lattice Resonance Band (Syncopated lattice only): 1/2-integer parametric resonance between *one* envelope mode and the lattice

Estimate mode frequencies for weak focusing from continuous focusing theory:

$$\sigma_B = \sigma_+ = \sqrt{2\sigma_0^2 + 2\sigma^2}$$

$$\sigma_Q = \sigma_- = \sqrt{\sigma_0^2 + 3\sigma^2}$$

This gives (measure phase advance in degrees):

Confluent Band:

$$(\sigma_+ + \sigma_-)/2 = 180^\circ$$

$$\Longrightarrow \quad \sqrt{2\sigma_0^2 + 2\sigma^2} + \sqrt{\sigma_0^2 + 3\sigma^2} = 360^\circ$$

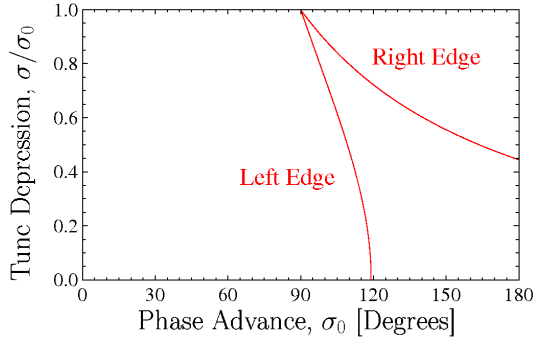
Lattice Resonance Band:

$$\sigma_+ = 180^\circ$$

$$\Longrightarrow \quad \sqrt{2\sigma_0^2 + 2\sigma^2} = 180^\circ$$

- Predictions poor due to inaccurate mode frequency estimates from continuous model
  - Predictions nearer to edge of band rather than center (expect resonance strongest at center)
- Cannot predict width of band
  - Important to characterize to avoid instability

To provide a rough guide on the location/width of the important FODO confluent instability band, many parametric runs were made and the instability region boundary was quantified through curve fitting:



Left Edge Boundary:

$$\sigma^2 + f(\eta)\sigma_0^2 = (90^\circ)^2[1 + f(\eta)]$$

$$f(\eta) = \frac{4}{3}$$

Right Edge Boundary:

$$\sigma + g(\eta)\sigma_0 = 90^\circ[1 + g(\eta)]$$

$$g(\eta) = \frac{1}{9}$$

- Negligible variation in quadrupole occupancy η is observed
- Formulas have a maximum error ~5 and ~2 degrees on left and right boundaries

Pure mode launching conditions for quadrupole focusing

Launching a pure breathing (B) or quadrupole (Q) mode in alternating gradient quadrupole focusing requires specific projections that generally require an eigenvalue/eigenvector analysis of symmetries to carry out

- See eignenvalue symmetries given in S6

Show example launch conditions for:

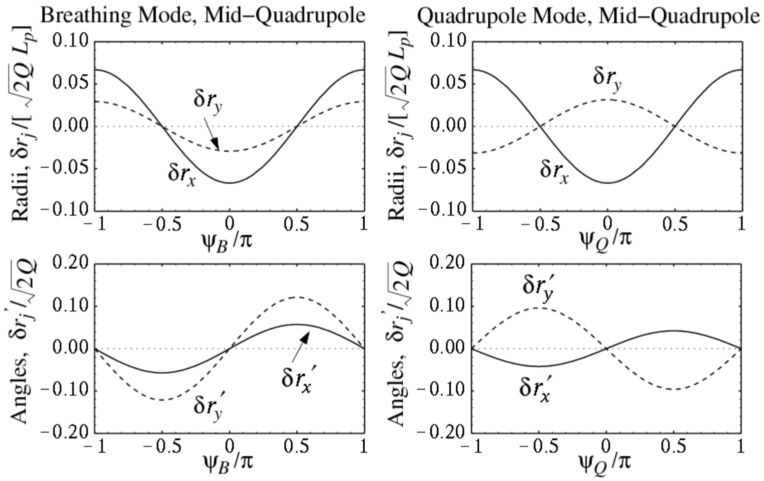
FODO Lattice

$$\eta = 0.6949$$

$$\sigma_0 = 80^\circ$$

$$\sigma/\sigma_0 = 0.2$$

Quadrupole Focusing – projections of perturbations on pure modes varies strongly with mode phase and the location in the lattice (FODO example)

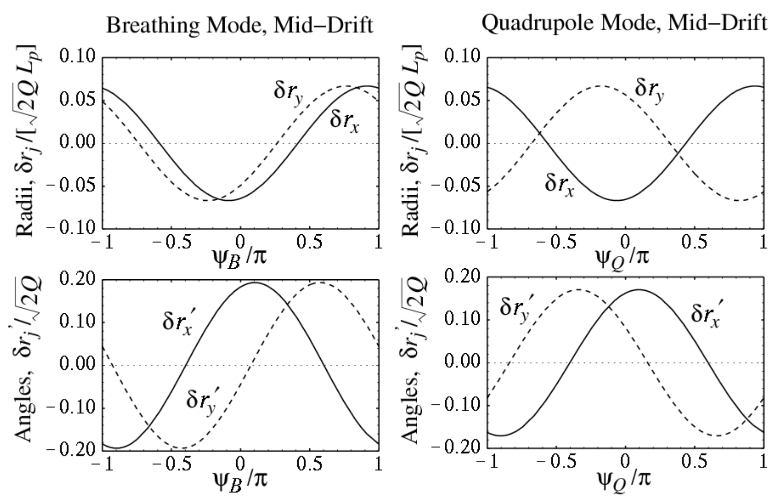


$\delta r_x \neq \delta r_y$

generally not exact breathing symmetry

$\delta r_x \neq -\delta r_y$

generally not exact quadrupole symmetry



$\delta r_x \neq \delta r_y$

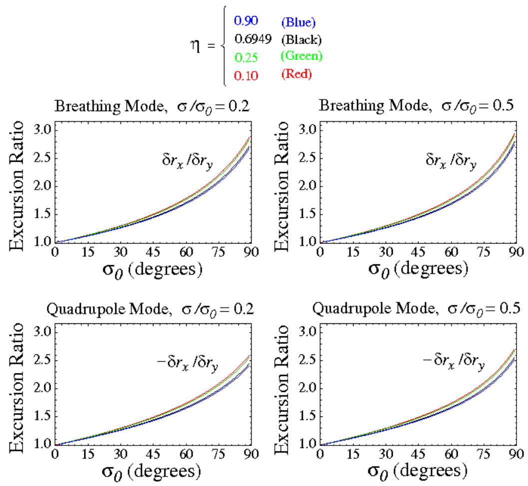
generally not exact breathing symmetry

$\delta r_x \neq -\delta r_y$

generally not exact quadrupole symmetry

As a further guide in pure mode launching, summarize FODO results for:

- Mid-axial x-focusing quadrupole with the additional choice $\delta r'_j = 0$
- Specify ratio of $\delta r_x/\delta r_y$ to launch pure mode
- Plot as function of σ_0 for $\sigma_0 < 90^\circ$
  - Results vary little with occupancy η or σ/σ_0



Comments:

- For quadrupole transport using the axisymmetric equilibrium projections on the breathing (+) mode and quadrupole (-) mode will *NOT* generally result in nearly pure mode projections:

$$\delta r_+ \equiv \frac{\delta r_x + \delta r_y}{2} \neq \text{Breathing Mode Projection}$$

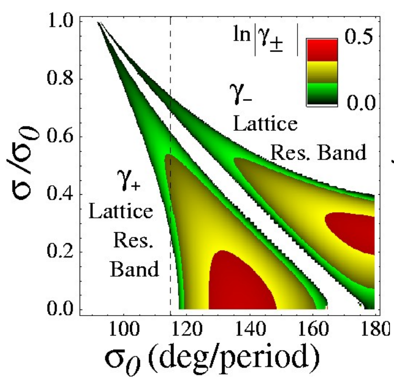
$$\delta r_- \equiv \frac{\delta r_x - \delta r_y}{2} \neq \text{Quadrupole Mode Projection}$$

  - Mistake can be commonly found in research papers and can confuse analysis of Suppposidly pure classes of envelope oscillations which are not.
  - Recall: reason denoted generalization of breathing mode with a subscript B and quadrupole mode with a subscript Q was an attempt to avoid confusion by overgeneralization
- Must solve for eigenvectors of 4x4 envelope transfer matrix through one lattice period calculated from the launch location in the lattice and analyze symmetries to determine proper projections (see S6)
- Normal mode coordinates can be found for the quadrupole and breathing modes in AG quadrupole focusing lattices through analysis of the eigenvectors but the expressions are typically complicated
  - Modes have underlying Courant-Snyder invariant but it will be a complicated

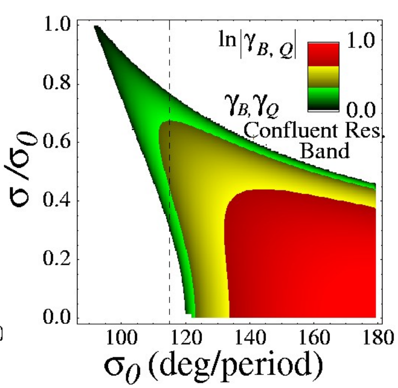
Summary: Envelope band instabilities and growth rates for periodic solenoidal and quadrupole doublet focusing lattices

Envelope Mode Instability Growth Rates

Solenoid (η = 0.25)



Quadrupole FODO (η = 0.70)



S9: Transport Limit Scaling Based on Envelope Models

See Handwritten Notes from 2008 USPAS

- Will convert to slides in future versions of the class

S10: Centroid and Envelope Descriptions via 1st Order Coupled Moment Equations

When constructing centroid and moment models, it can be efficient to simply write moments, differentiate them, and then apply the equation of motion. Generally, this results in lower order moments coupling to higher order ones and an infinite chain of equations. But the hierarchy can be truncated by:

- Assuming a fixed functional form of the distribution in terms of moments
- Neglecting coupling to higher order terms

Resulting first order moment equations can be expressed in terms of a closed set of moments and advanced in s or t using simple (ODE based) numerical codes. This approach can prove simpler to include effects where invariants are not easily extracted to reduce the form of the equations (as when solving the KV envelope equations in the usual form).

Examples of effects that might be more readily analyzed:

- Skew coupling in quadrupoles
- Chromatic effects in final focus
- Dispersion in bends

See: references at end of notes
J.J. Barnard, lecture on
Heavy-Ion Fusion and Final Focusing

Resulting 1st order form of coupled moment equations:

$$\frac{d}{ds}\mathbf{M} = \mathbf{F}(\mathbf{M})$$

$\mathbf{M}$ = vector of moments, and their s derivatives, generally infinite
$\mathbf{F}$ = vector function of $\mathbf{M}$, generally nonlinear

- System advanced from a specified initial condition (initial value of $\mathbf{M}$)

Transverse moment definition:

$$\langle \cdots \rangle_\perp \equiv \frac{\int d^2x_\perp \int d^2x'_\perp \cdots f_\perp}{\int d^2x_\perp \int d^2x'_\perp f_\perp}$$

Can be generalized if other variables such as off momentum are included in f

Differentiate moments and apply equations of motion:

$$\frac{d}{ds}\langle \cdots \rangle_\perp \equiv \frac{\int d^2x_\perp \int d^2x'_\perp \left[\frac{d}{ds}\cdots\right] f_\perp}{\int d^2x_\perp \int d^2x'_\perp f_\perp}$$

+ apply equations of motion to simplify $\frac{d}{ds}\cdots$

When simplifying the results, if the distribution form is frozen in terms of moments (Example: assume uniform density elliptical beam) then we use constructs like:

$$n = \int d^2x'_\perp \; f_\perp \;\; = n(\mathbf{M})$$

to simplify the resulting equations and express the RHS in terms of elements of $\mathbf{M}$

1st order moments:

$\mathbf{X}_\perp = \langle \mathbf{x}_\perp \rangle_\perp$ Centroid coordinate
$\mathbf{X}'_\perp = \langle \mathbf{x}'_\perp \rangle_\perp$ Centroid angle

+ possible others if more variables. Example

$\Delta = \langle \frac{\delta p_s}{p_s} \rangle = \langle \delta \rangle$ Centroid off-momentum

$\vdots$ $\vdots$

2nd order moments:

It is typically convenient to subtract centroid from higher-order moments

$$\tilde{x} \equiv x - X \qquad \tilde{x}' \equiv x' - X'$$
$$\tilde{y} \equiv y - Y \qquad \tilde{y}' \equiv y' - Y'$$
$$\tilde{\delta} \equiv \delta - \Delta$$

x-moments	*y*-moments	*x*-*y* cross moments	dispersive moments
$\langle \tilde{x}^2 \rangle_\perp$	$\langle \tilde{y}^2 \rangle_\perp$	$\langle \tilde{x}\tilde{y} \rangle_\perp$	$\langle \tilde{x}\tilde{\delta} \rangle$, $\langle \tilde{y}\tilde{\delta} \rangle$
$\langle \tilde{x}\tilde{x}' \rangle_\perp$	$\langle \tilde{y}\tilde{y}' \rangle_\perp$	$\langle \tilde{x}'\tilde{y} \rangle_\perp$, $\langle \tilde{x}\tilde{y}' \rangle_\perp$	$\langle \tilde{x}'\tilde{\delta} \rangle$, $\langle \tilde{y}'\tilde{\delta} \rangle$
$\langle \tilde{x}'^2 \rangle_\perp$	$\langle \tilde{y}'^2 \rangle_\perp$	$\langle \tilde{x}'\tilde{y}' \rangle_\perp$	$\langle \tilde{\delta}^2 \rangle$

3rd order moments: Analogous to 2nd order case, but more for each order

$$\langle \tilde{x}^3 \rangle_\perp, \; \langle \tilde{x}^2\tilde{y} \rangle_\perp, \; \cdots$$

Many quantities of physical interest are expressed in transport can then be expressed in terms of moments calculated when the equations are numerically advanced in *s* and their evolutions plotted to understand behavior

- Many quantities of physical interest are expressible in terms of 1st and 2nd order moments

Example moments often projected:

Statistical beam size:
(rms edge measure)

$$r_x = 2\langle \tilde{x}^2 \rangle_\perp^{1/2}$$
$$r_y = 2\langle \tilde{y}^2 \rangle_\perp^{1/2}$$

Statistical emittances:
(rms edge measure)

$$\varepsilon_x = 4\left[\langle \tilde{x}^2 \rangle_\perp \langle \tilde{x}'^2 \rangle_\perp - \langle \tilde{x}\tilde{x}' \rangle_\perp^2\right]^{1/2}$$
$$\varepsilon_y = 4\left[\langle \tilde{y}^2 \rangle_\perp \langle \tilde{y}'^2 \rangle_\perp - \langle \tilde{y}\tilde{y}' \rangle_\perp^2\right]^{1/2}$$

Kinetic longitudinal temperature:
(rms measure)

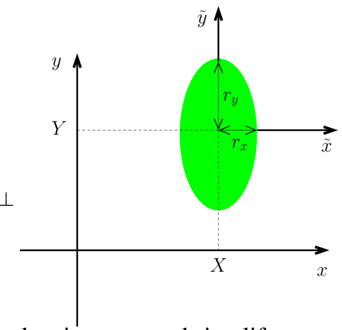
$$T_s = \text{const} \times \langle \tilde{\delta}^2 \rangle$$

Illustrate approach with the familiar KV model

Truncation assumption: unbunched uniform density elliptical beam in free space

- $\delta = 0$, no axial velocity spread
- All cross moments zero, i.e. $\langle \tilde{x}\tilde{y} \rangle_\perp = 0$

$$\frac{d}{ds}\langle x \rangle_\perp = \langle x' \rangle_\perp \qquad \frac{d}{ds}\langle x^2 \rangle_\perp = 2\langle x x' \rangle_\perp$$
$$\frac{d}{ds}\langle x' \rangle_\perp = \langle x'' \rangle_\perp \qquad \frac{d}{ds}\langle x'^2 \rangle_\perp = 2\langle x' x'' \rangle_\perp$$
$$\vdots \qquad\qquad \vdots$$



Use particle equations of motion within beam, neglect images, and simplify

- Apply equations in S2 with $\mathbf{E}_\perp^i = 0$

$$x'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}x' + \kappa_x x - \frac{2Q}{(r_x + r_y)r_x}(x - \langle x \rangle_\perp) = 0$$
$$y'' + \frac{(\gamma_b\beta_b)'}{(\gamma_b\beta_b)}y' + \kappa_y y - \frac{2Q}{(r_x + r_y)r_y}(y - \langle y \rangle_\perp) = 0$$

Resulting system of 1st and 2nd order moments

1st order moments:

$$\frac{d}{ds}\begin{bmatrix} \langle x \rangle_\perp \\ \langle x' \rangle_\perp \\ \langle y \rangle_\perp \\ \langle y' \rangle_\perp \end{bmatrix} = \begin{bmatrix} \langle x' \rangle_\perp \\ -\kappa_x(s)\langle x \rangle_\perp \\ \langle y' \rangle_\perp \\ -\kappa_y(s)\langle y \rangle_\perp \end{bmatrix}$$

2nd order moments:

$$\frac{d}{ds}\begin{bmatrix} \langle \tilde{x}^2 \rangle_\perp \\ \langle \tilde{x}\tilde{x}' \rangle_\perp \\ \langle \tilde{x}'^2 \rangle_\perp \\ \langle \tilde{y}^2 \rangle_\perp \\ \langle \tilde{y}\tilde{y}' \rangle_\perp \\ \langle \tilde{y}'^2 \rangle_\perp \end{bmatrix} = \begin{bmatrix} 2\langle \tilde{x}\tilde{x}' \rangle_\perp \\ \langle \tilde{x}'^2 \rangle_\perp - \kappa_x(s)\langle \tilde{x}^2 \rangle_\perp + \frac{Q\langle \tilde{x}^2 \rangle_\perp^{1/2}}{2[\langle \tilde{x}^2 \rangle_\perp^{1/2} + \langle \tilde{y}^2 \rangle_\perp^{1/2}]} \\ -2\kappa_x(s)\langle \tilde{x}\tilde{x}' \rangle_\perp + \frac{Q\langle \tilde{x}\tilde{x}' \rangle_\perp}{\langle \tilde{x}^2 \rangle_\perp^{1/2}[\langle \tilde{x}^2 \rangle_\perp^{1/2} + \langle \tilde{y}^2 \rangle_\perp^{1/2}]} \\ 2\langle \tilde{y}\tilde{y}' \rangle_\perp \\ \langle \tilde{y}'^2 \rangle_\perp - \kappa_y(s)\langle \tilde{y}^2 \rangle_\perp + \frac{Q\langle \tilde{y}^2 \rangle_\perp^{1/2}}{2[\langle \tilde{x}^2 \rangle_\perp^{1/2} + \langle \tilde{y}^2 \rangle_\perp^{1/2}]} \\ -2\kappa_y(s)\langle \tilde{y}\tilde{y}' \rangle_\perp + \frac{Q\langle \tilde{y}\tilde{y}' \rangle_\perp}{\langle \tilde{y}^2 \rangle_\perp^{1/2}[\langle \tilde{x}^2 \rangle_\perp^{1/2} + \langle \tilde{y}^2 \rangle_\perp^{1/2}]} \end{bmatrix}$$

- Express 1st and 2nd order moments separately in this case since uncoupled
- Form truncates due to frozen distribution form: all moments on LHS on RHS
- Integrate from initial moments values of *s* and project out desired quantities

Using 2nd order moment equations we can show that

$$\frac{d}{ds}\varepsilon_x^2 = 0 = \frac{d}{ds}\varepsilon_y^2$$

$$\Longrightarrow \quad \begin{aligned} \varepsilon_x^2 &= 16\left[\langle x^2 \rangle_\perp \langle x'^2 \rangle_\perp - \langle x x' \rangle_\perp^2\right] = \text{const} \\ \varepsilon_y^2 &= 16\left[\langle y^2 \rangle_\perp \langle y'^2 \rangle_\perp - \langle y y' \rangle_\perp^2\right] = \text{const} \end{aligned}$$

Using this, the 2nd order moment equations can be equivalently expressed in the standard KV envelope form:

$$\frac{dr_x}{ds} = r_x' \; ; \quad \frac{d}{ds}r_x' + \kappa_x r_x - \frac{2Q}{r_x + r_y} - \frac{\varepsilon_x^2}{r_x^3} = 0$$
$$\frac{dr_y}{ds} = r_y' \; ; \quad \frac{d}{ds}r_y' + \kappa_y r_y - \frac{2Q}{r_x + r_y} - \frac{\varepsilon_y^2}{r_y^3} = 0$$

- Moment form fully consistent with usual KV model .... as it must be
- Moment form generally easier to put in additional effects that would violate the usual emittance invariants

Relative advantages of the use of coupled matrix form versus reduced equations can depend on the problem/situation

Coupled Matrix Equations

$$\frac{d}{ds}\mathbf{M} = \mathbf{F}(\mathbf{M})$$

M = Moment Vector
F = Force Vector

- Easy to formulate
  - Straightforward to incorporate additional effects
- Natural fit to numerical routine
  - Easy to numerically code/solve

Reduced Equations

$$X'' + \kappa_x X = 0$$

$$r_x'' + \kappa_x r_x - \frac{2Q}{r_x + r_y} - \frac{\varepsilon_x^2}{r_x^3} = 0$$

etc.

Reduction based on identifying invariants such as $\varepsilon_x^2 = 16\left[\langle\tilde{x}^2\rangle_\perp \langle\tilde{x}'^2\rangle_\perp - \langle\tilde{x}\tilde{x}'\rangle_\perp^2\right]$ helps understand solutions

- Compact expressions can help analytical understanding

These notes will be corrected and expanded for reference and future editions of US Particle Accelerator School and University of California at Berkeley courses:
"Beam Physics with Intense Space Charge"
"Interaction of Intense Charged Particle Beams with Electric and Magnetic Fields"
by J.J. Barnard and S.M. Lund

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